

# Recent improvement of operational models

## I. Modeling of in-ground effects

## II. Modeling of atmospheric stratification and turbulence effects

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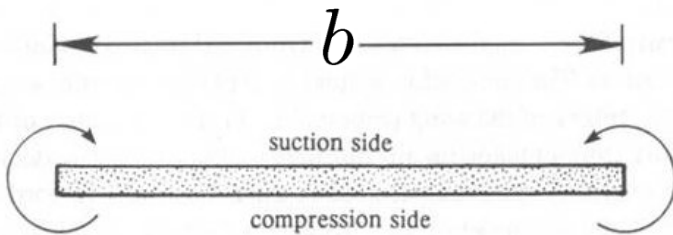
# Aircraft wake vortices

$$L = \rho U_{\infty} b_0 \Gamma_0 = M g$$

$$\Gamma_0 = \frac{M g}{\rho U_{\infty} b_0} \quad \text{circulation of WV}$$

$$V_0 = \frac{\Gamma_0}{2\pi b_0} \quad \text{sink velocity of WV}$$

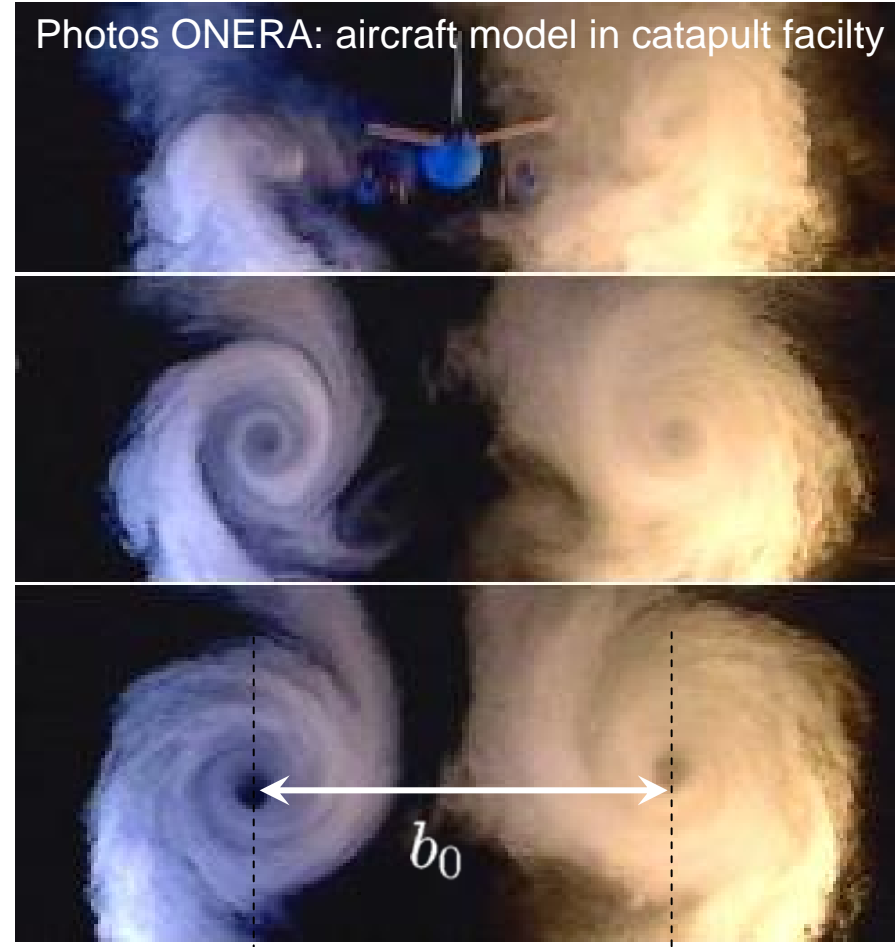
$$t_0 = \frac{b_0}{V_0} \quad \text{reference time}$$



$$b_0 = 0.66 \dots 0.92 b$$

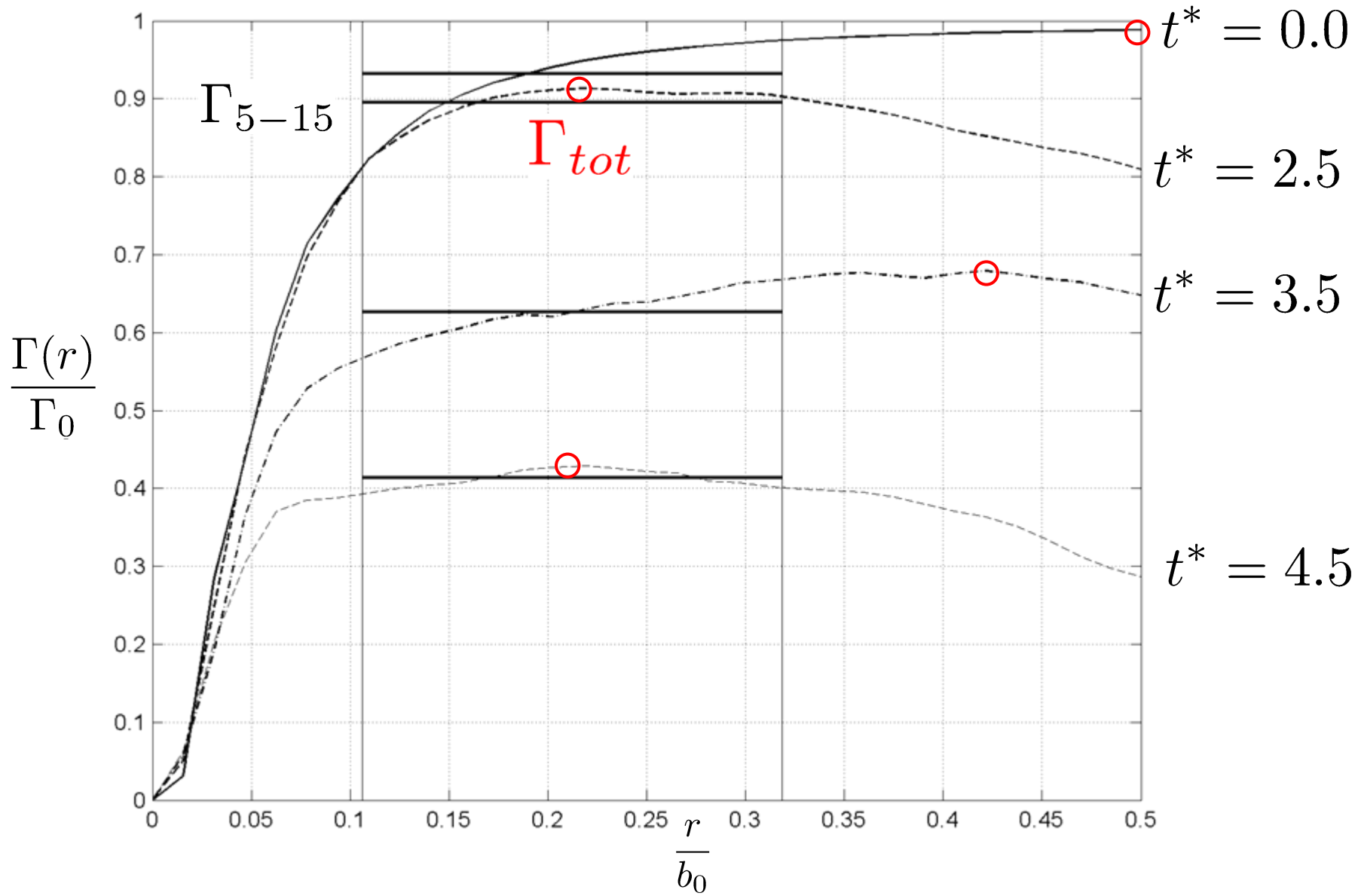
depending on a/c and on wing configuration

Photos ONERA: aircraft model in catapult facility



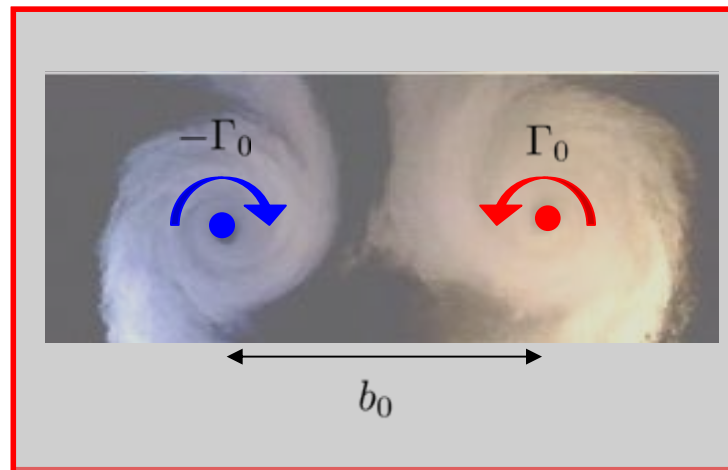
# Circulation definition

$$\Gamma(r, t) = \int_0^{2\pi} \int_0^r \omega_x(r', \theta, t) r' dr' d\theta$$



# The Deterministic wake Vortex Model (DVM)

- The DVM forecasts, in real-time, the WV behavior (transport and decay) in one computational gate, using simplified physical models.
- Each wake vortex is represented by a vortex particle (or by a set of vortex particles)
- The following effects are modeled for transport and decay:
  - Biot-Savart transport,
  - influence of the wind (wind transport, wind shear),
  - influence of atmospheric turbulence and stratification,
  - influence of the ground proximity.



# Ground proximity modeling

- Use of secondary vortex particles generated close to the ground
- Interaction with the primary vortices induces :
  - the rebound,
  - an enhanced decay
- Model much improved in the framework of the FAR-Wake project (WP3- Wake Vortices In-Ground Effect)
- Recent improvements of the model to be computationally more efficient

# Near Ground Effect Modeling

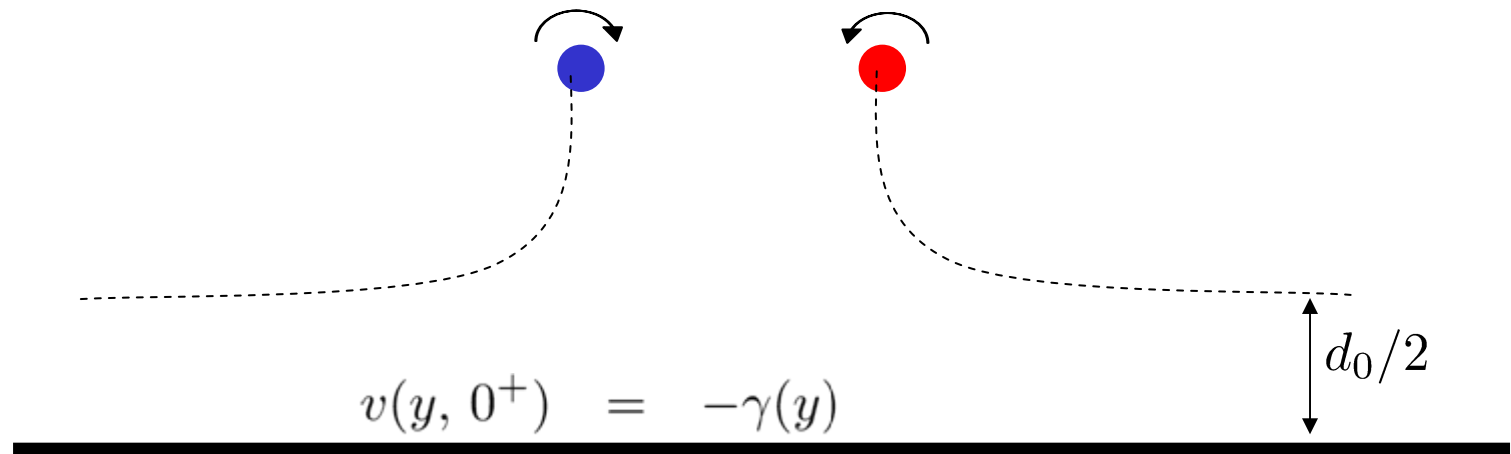
- Use of images vortex particles below the ground
- This is like computing  $\gamma(y)$  (which would require solving a boundary integral equation)
- Correct prediction of the inviscid transport



$$d\Gamma(y) = \gamma(y) dy \quad \begin{array}{l} v(y, 0^+) = -\gamma(y) \\ v(y, 0) = 0 \end{array}$$

# Near Ground Effect Modeling

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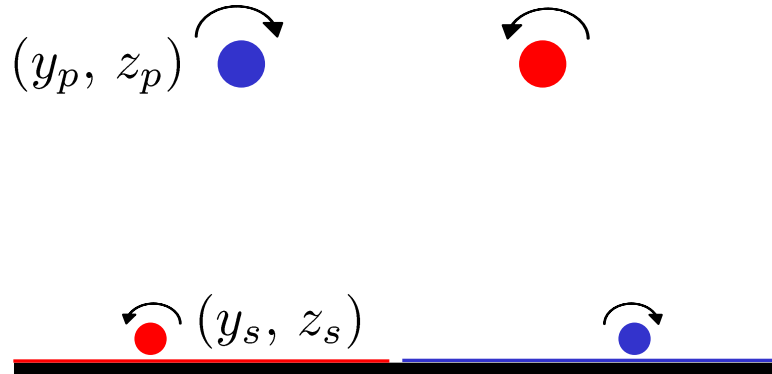
with 
$$\frac{d_0}{2} = \frac{b_0}{\sqrt{\left(\frac{b_0}{h_0}\right)^2 + 4}}$$



# In-Ground Effect Modeling

Modeling of the separation of the boundary layer generated at the ground:

- At each time-step, generation of a new secondary particle when  $z_p \leq \alpha_s \frac{d_0}{2}$   
(typically  $\alpha_s = 2.4$ )

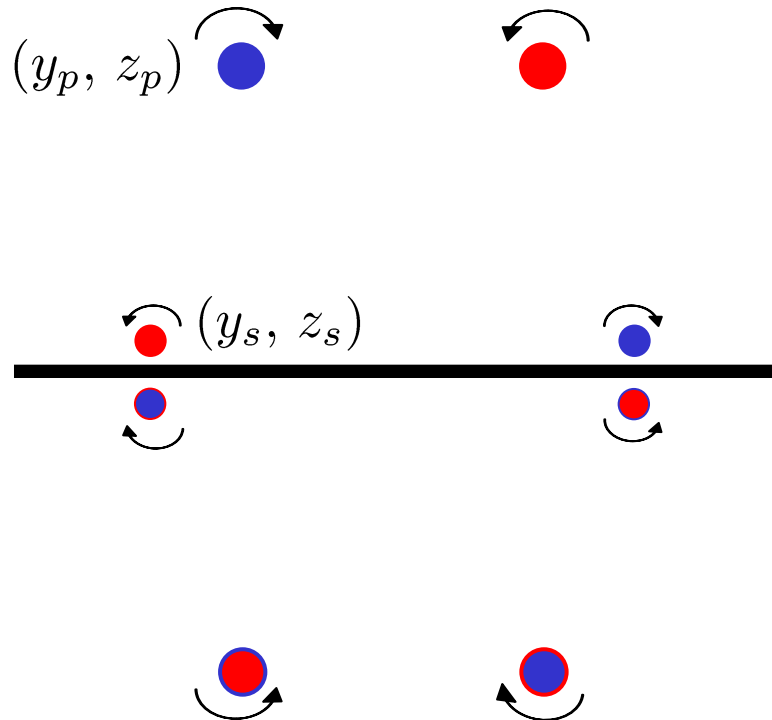




# In-Ground Effect Modeling

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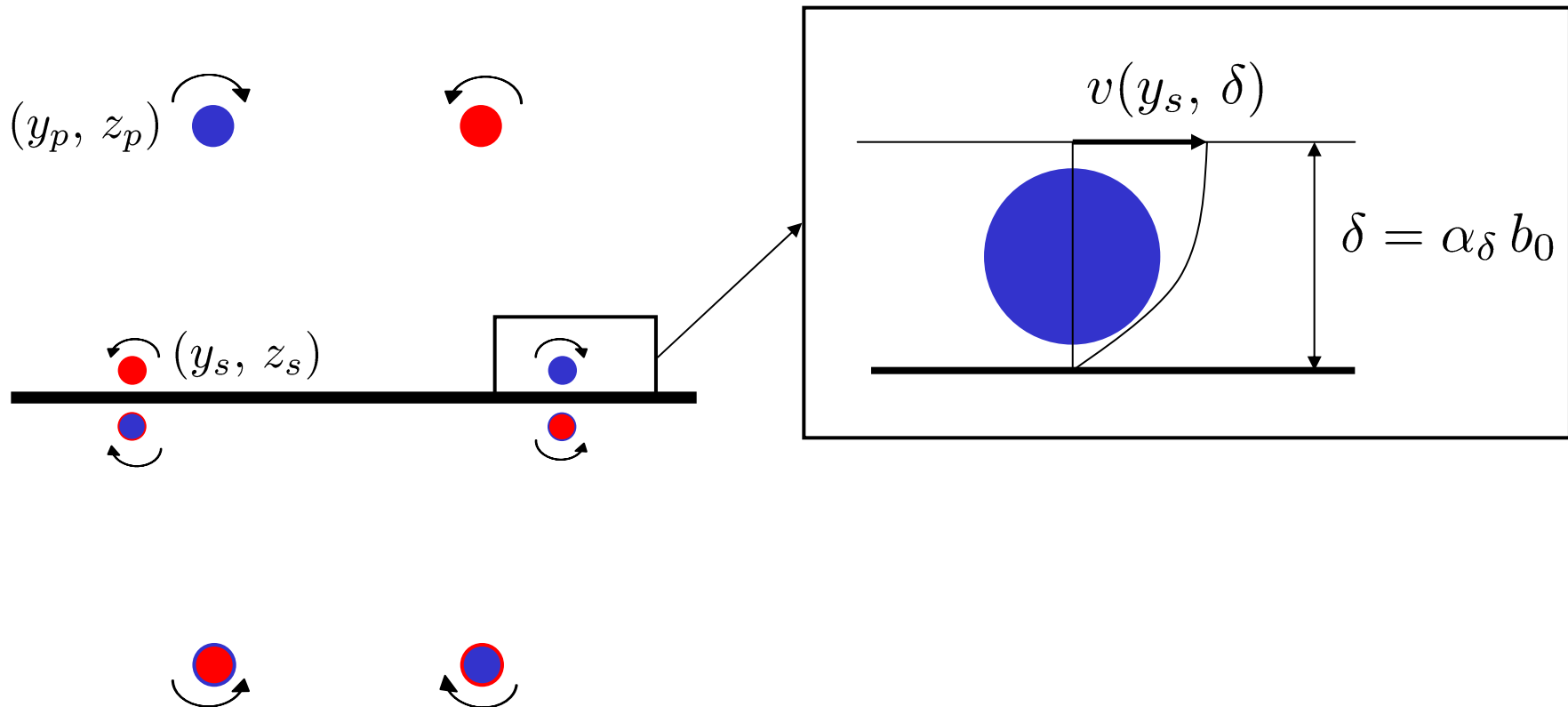


# In-Ground Effect Modeling

Secondary particles characteristics are function of the whole vortex system and of

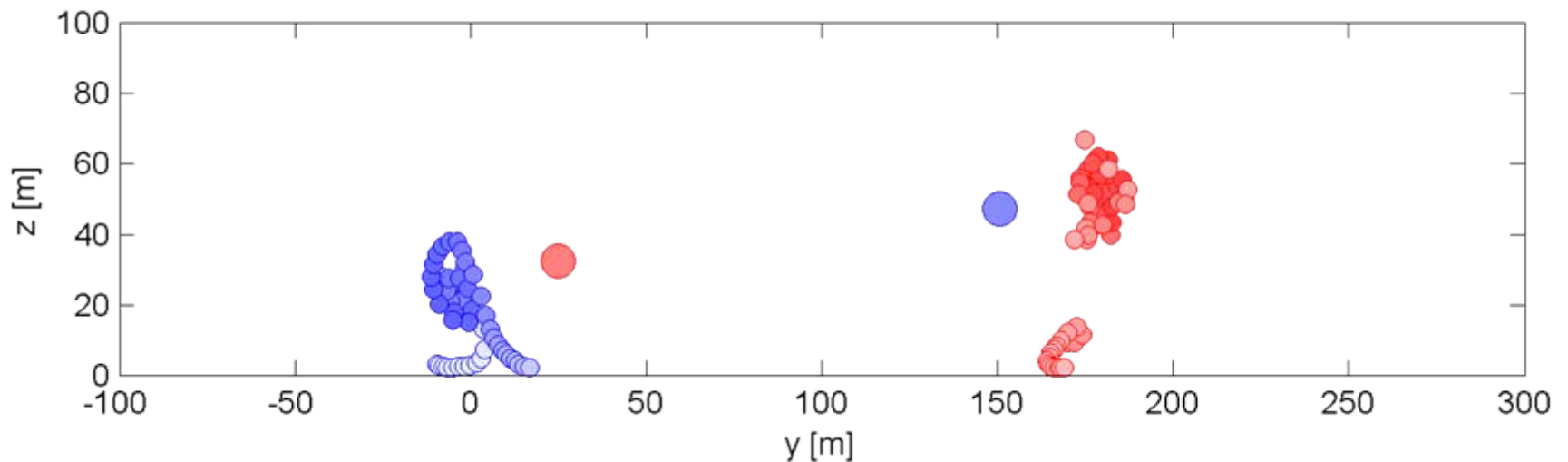
the x-wind:

$$\left\{ \begin{array}{lcl} v(y_s, 0) & = & \alpha_v \max(v(y, 0)) \\ \Gamma_s & = & \pm \alpha_\Gamma \frac{v^2(y_s, \delta)}{\frac{\delta}{2}} \Delta t \\ z_s & = & \end{array} \right. \quad \left\{ \begin{array}{lcl} \alpha_v & = & 0.82 \\ \alpha_\Gamma & = & 1.5 \\ \alpha_\delta & = & 0.08 \end{array} \right.$$



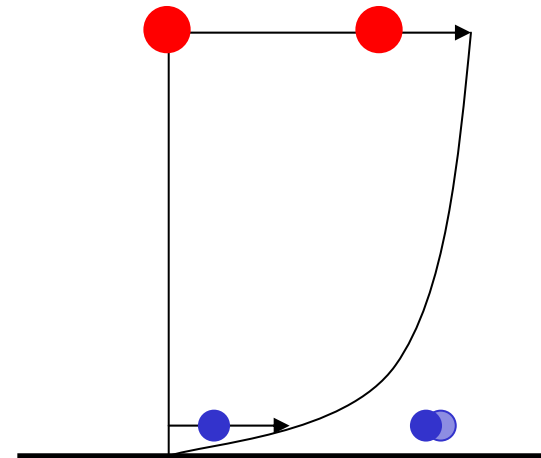
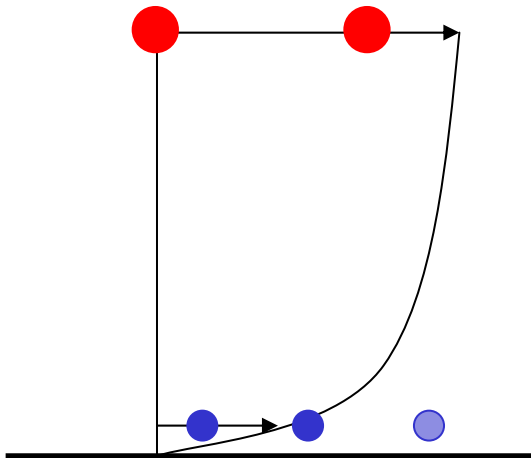
# Dynamic separation of the secondary particles

- Using Biot-Savart and wind transport:
  - The particles dynamically separate from the ground
  - They interact with the primary vortices and induce the rebound
- As the position and strength of the secondary vortex particles also depend on the x-wind, it also reproduces the asymmetric rebound



# Special treatment of wind transport for particles very close to the ground

- Due to the wind profile, the wind speed very close to the ground can be a lot smaller than aloft.
- In case of strong crosswind, the secondary vortex particles would not “follow” the primary ones.
- The secondary vortex particles have then “no time” to separate.
- Solution: the wind speed, used for the secondary particles located very close to the ground (typically below 10 m), is evaluated at the primary vortex altitude.



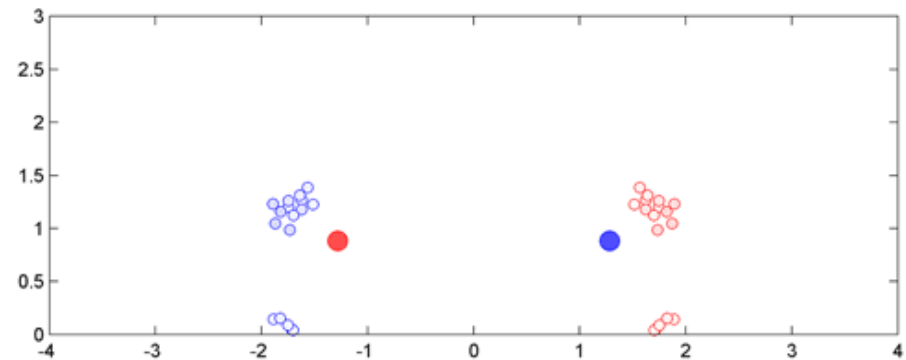
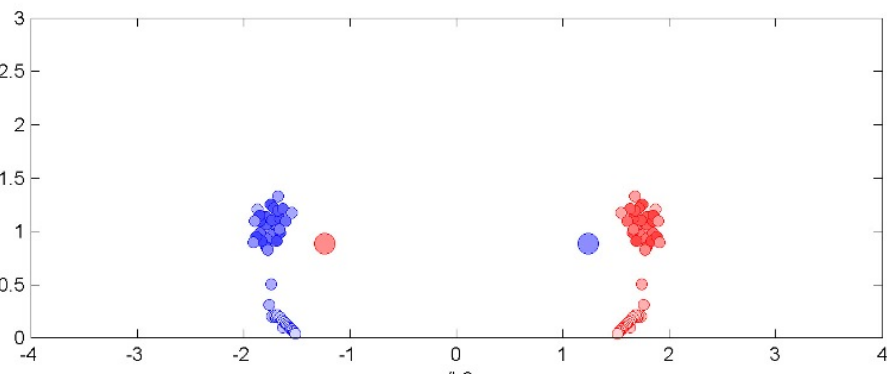
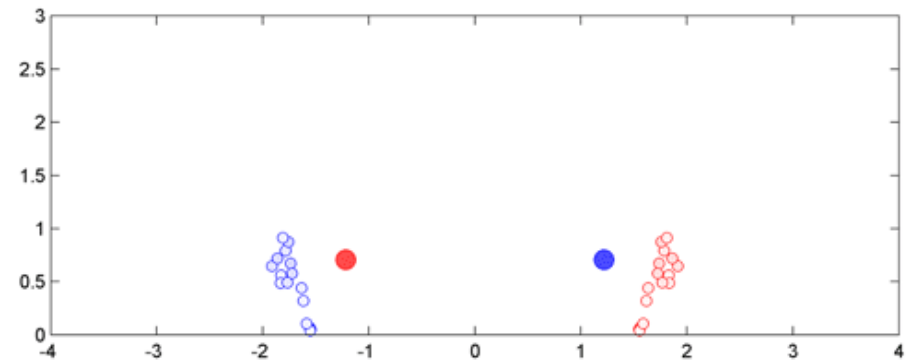
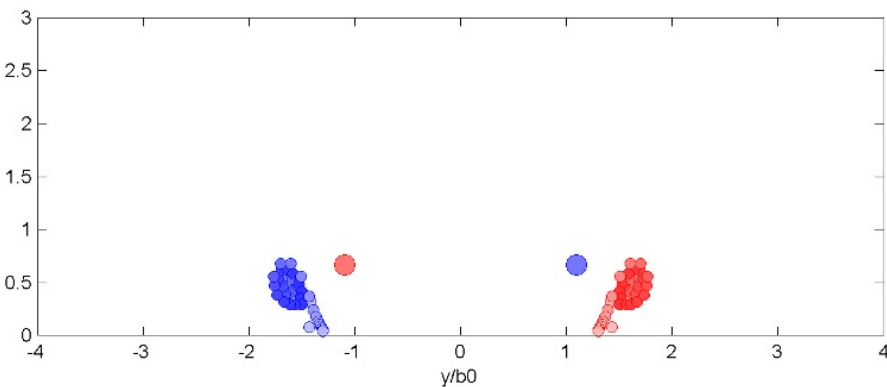
# Enhanced IGE decay model

- Model based on a “Particle Strength Exchange” (PSE) approach
- Exponential decay of both the primary and secondary particles
- The decay rate of the primary (resp. secondary) particles depends on:
  - the total secondary (resp. primary) particle circulation
  - the minimum altitude reached by the primary particles
- System applied separately to the port and starboard vortices:

$$\begin{aligned}\frac{d\Gamma_p}{dt} &= -C_{IGE} \frac{|\sum_s \Gamma_s|}{(z_{min})^2} \Gamma_p, \\ \frac{d\Gamma_s}{dt} &= -C_{IGE} \frac{|\sum_p \Gamma_p|}{(z_{min})^2} \Gamma_s.\end{aligned}\quad C_{IGE} = 0.017$$

# Agglomeration of secondary particles

- The IGE model requires the use of several secondary vortex particles.
- Solution: agglomeration of particles
- Once a structure made of several particles reaches:
  - a sufficient size, and/or
  - a sufficient circulation,
 it is replaced by one equivalent particle



# Reorganization of the secondary particles

- After the first rebound, coherent structures of secondary vorticity are no longer observed around the primary vortex.
- This is a 3-D effect => not dynamically captured by the 2-D vortex particle method.
- Model: once the primary vortex has rebound and tends to sink again, we reorganize the secondary vortex particles in a “halo” around the primary vortex.

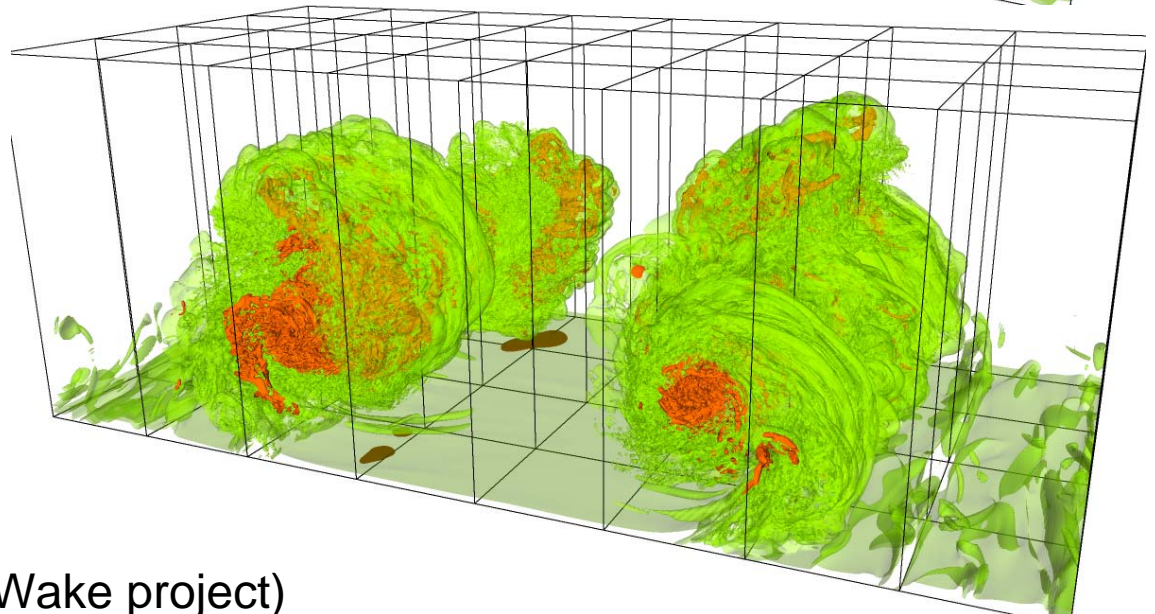
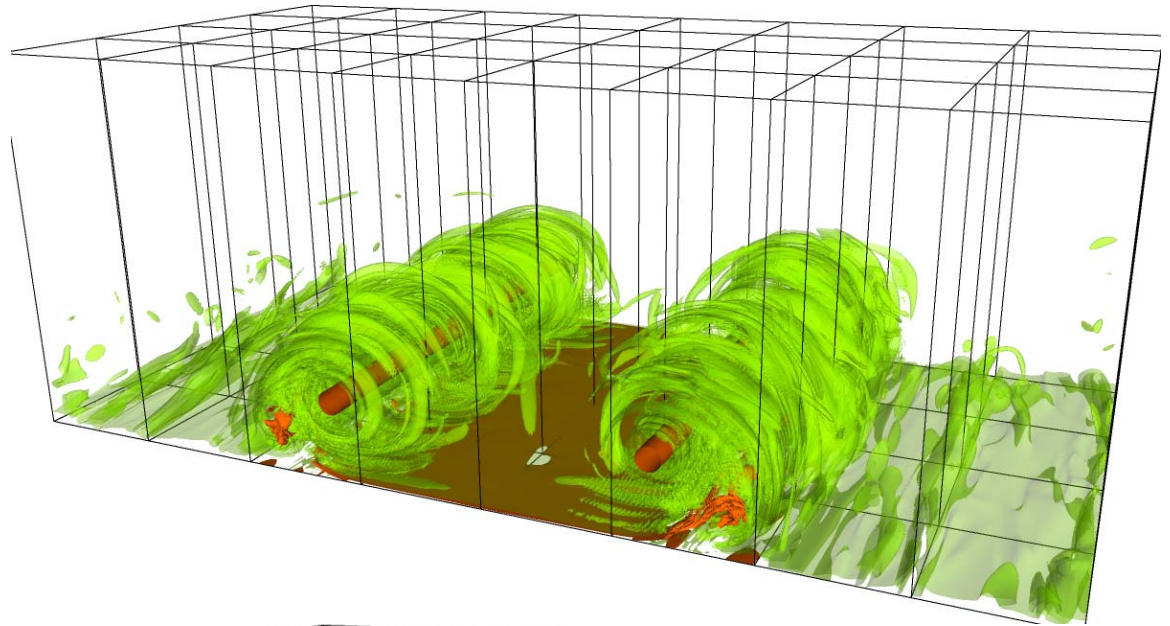


Source: UCL (FAR-Wake project)

# LES of a 2VS IGE with headwind

$$h_0 = b_0$$

$$V_{\text{wind}}(h_0) = V_0$$

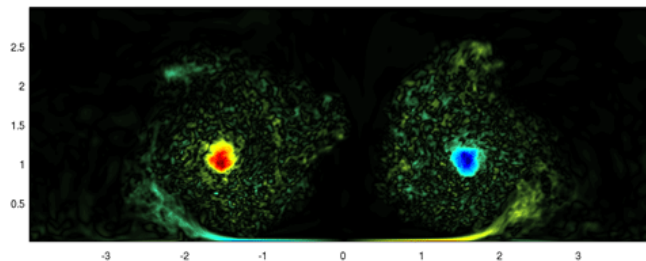
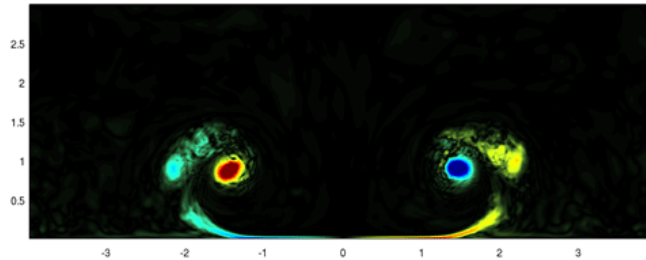
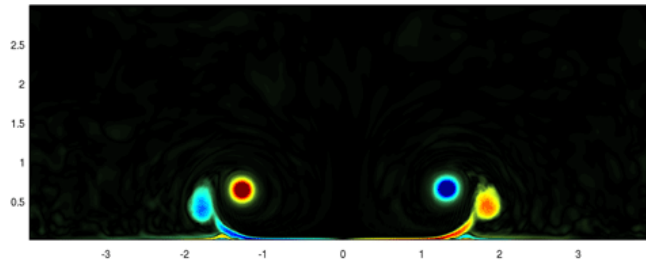
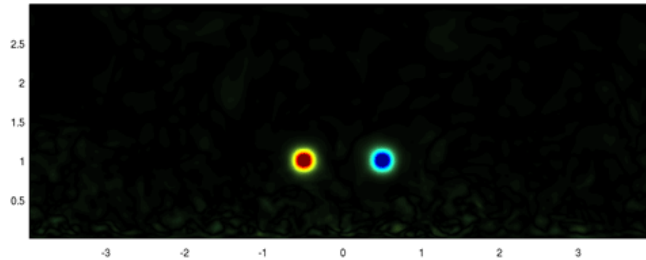


Source: Cenaero (FAR-Wake project)

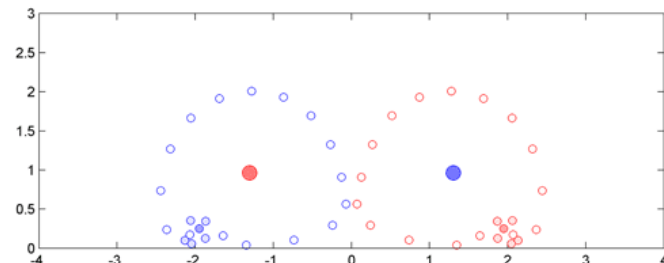
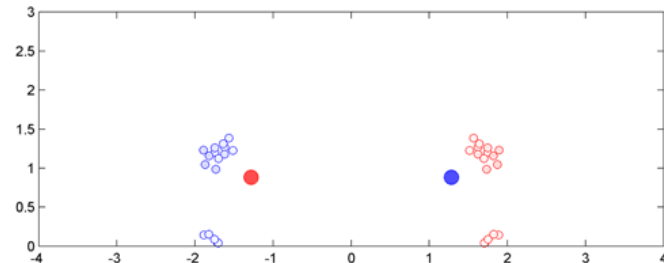
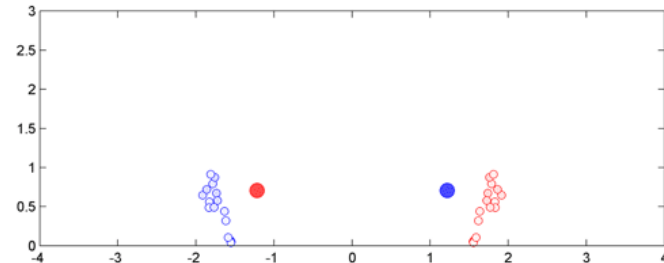
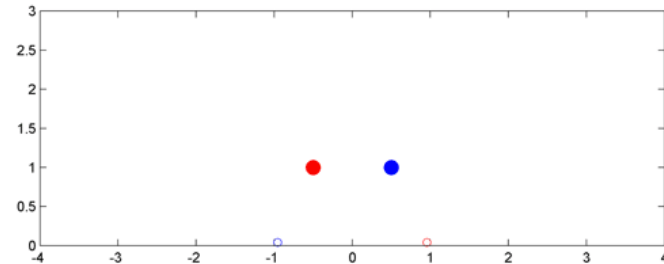


# Comparison between LES and DVM results: - case IGE with headwind

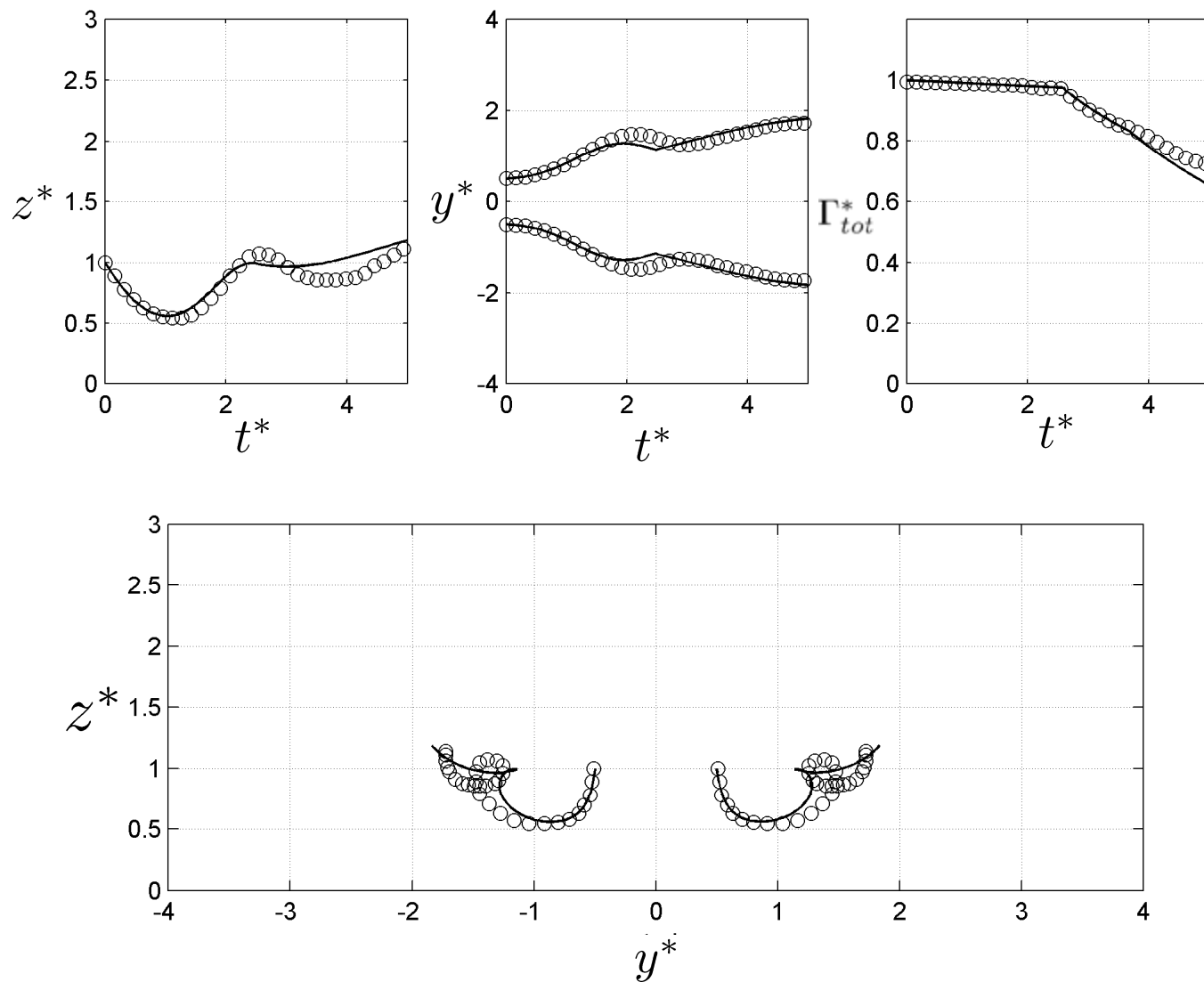
*Mean axial vorticity field as computed by a LES*



*Particles used in the DVM to model WV IGE*



# Comparison between LES and DVM results: - case IGE with headwind

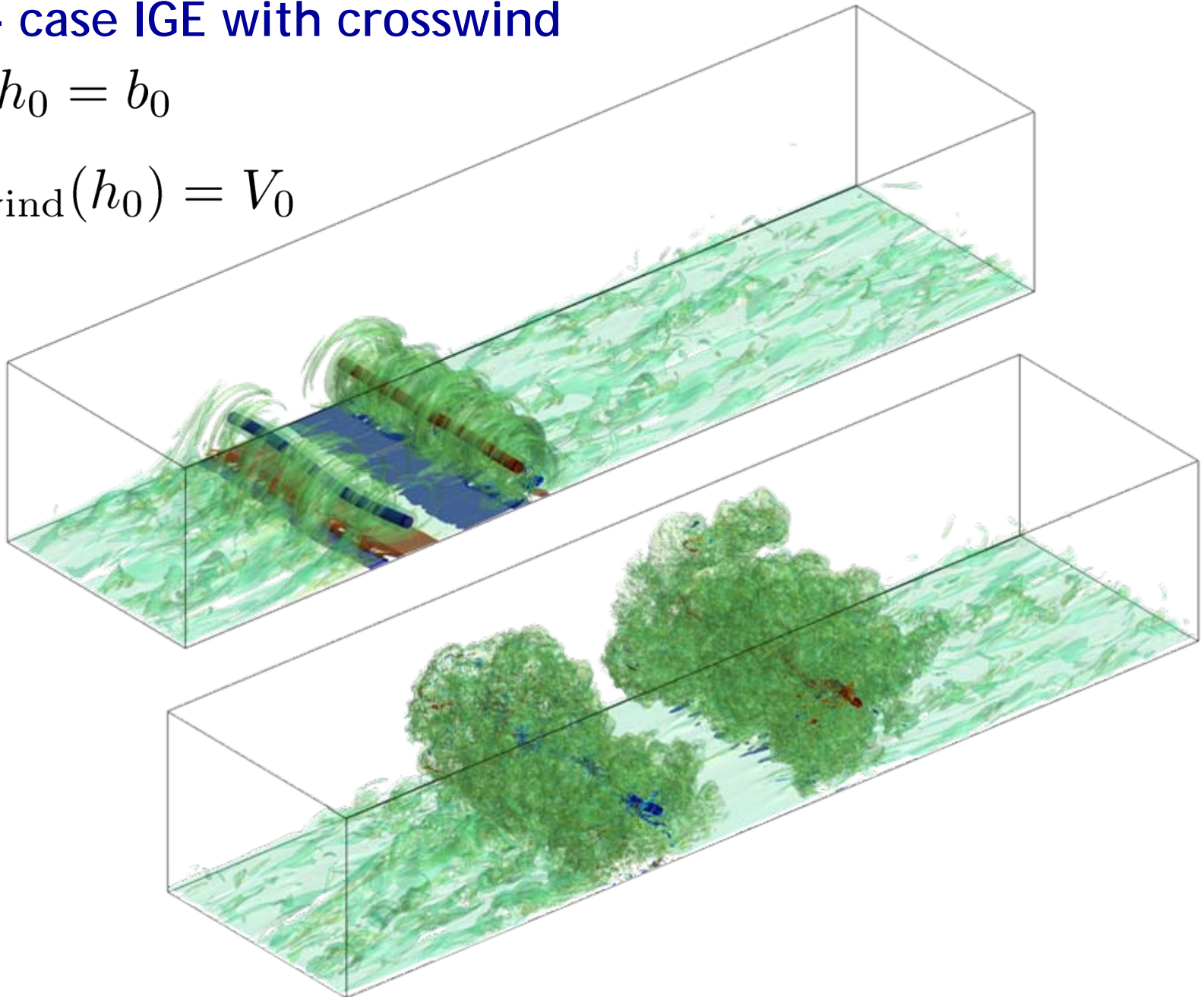


# Comparison between LES and DVM results:

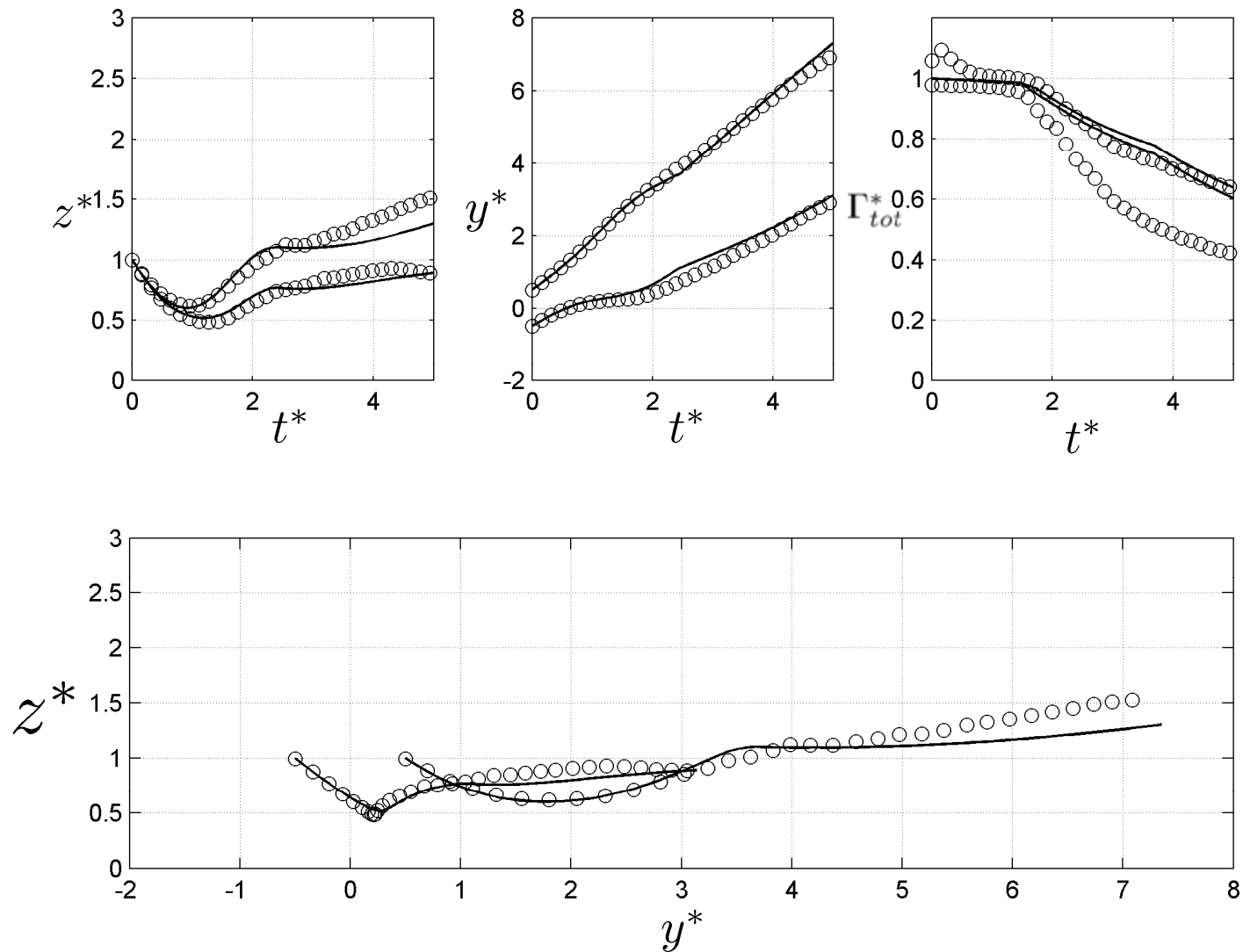
- case IGE with crosswind

$$h_0 = b_0$$

$$V_{\text{wind}}(h_0) = V_0$$

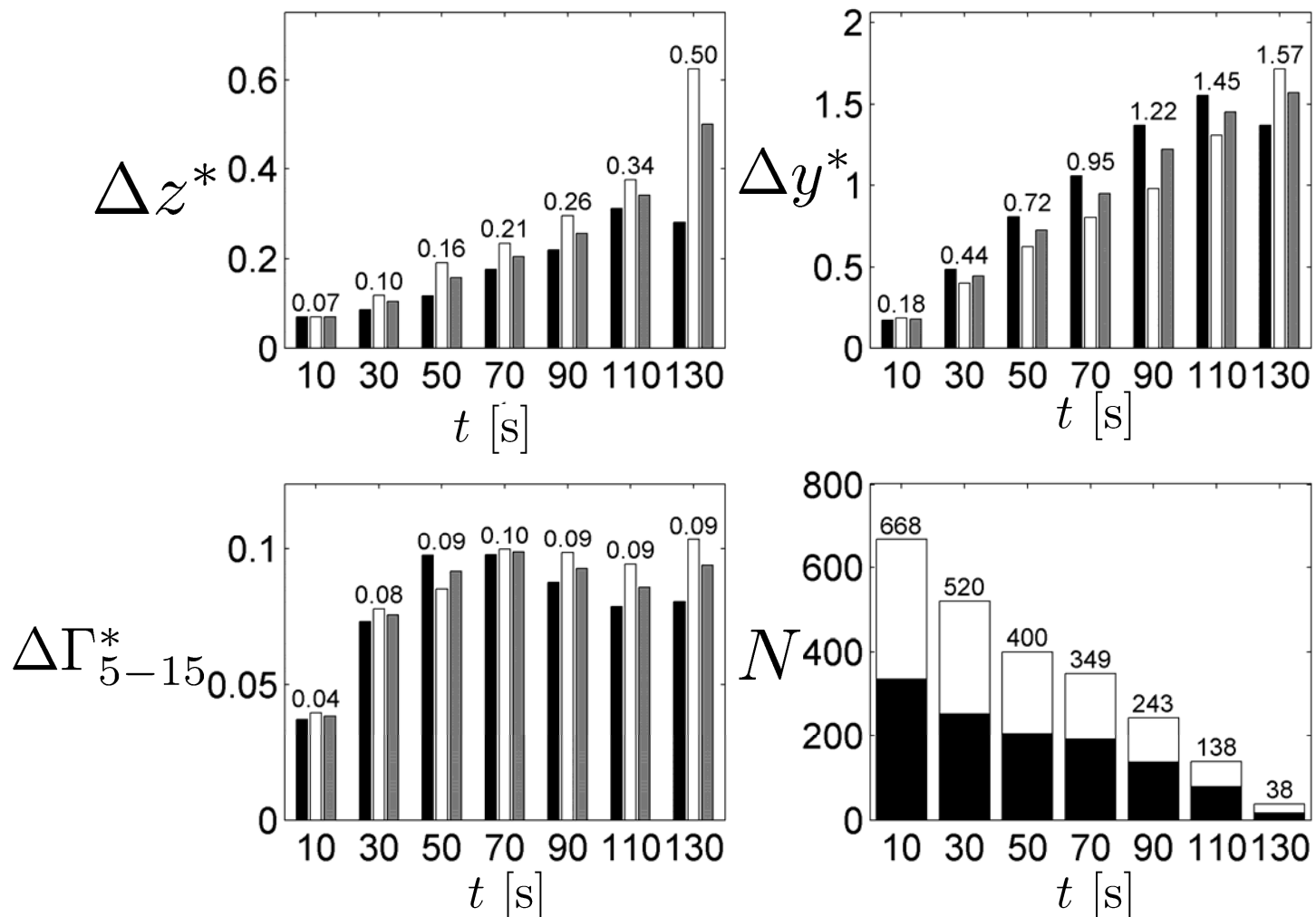


# Comparison between LES and DVM results: - case IGE with crosswind



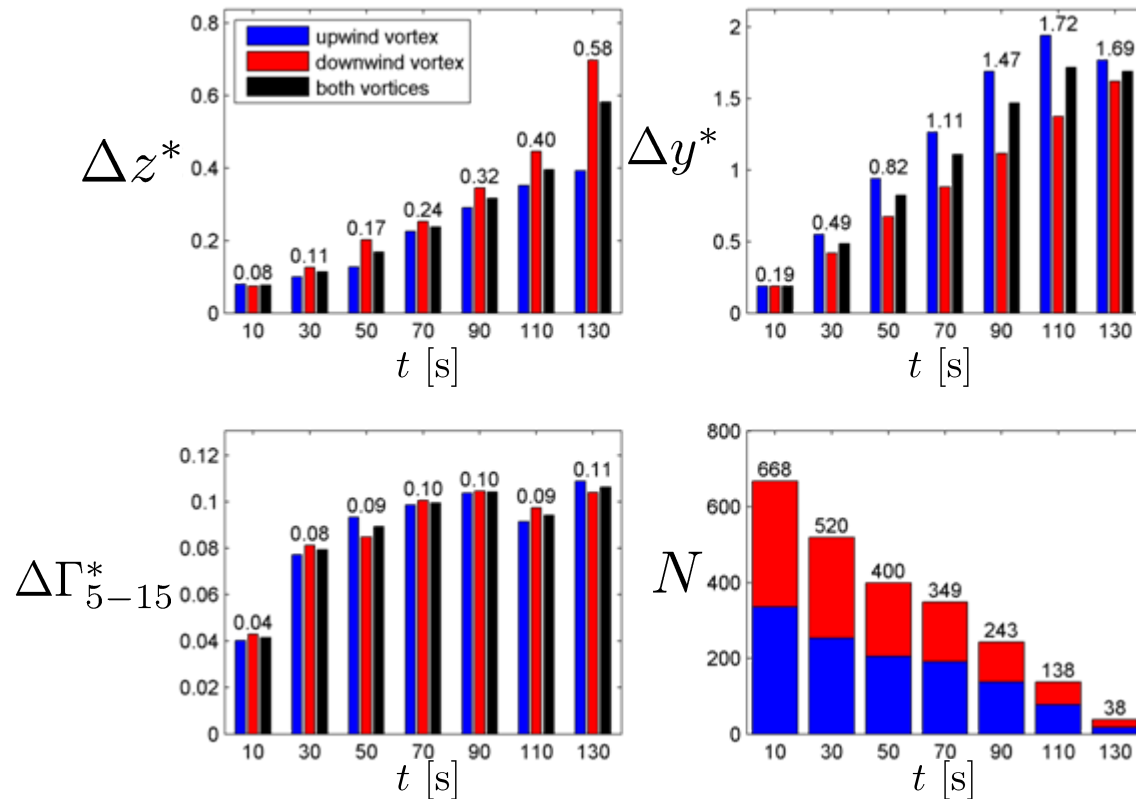
# WakeFRA 2004 database

- LIDAR measurement campaign of WV generated by heavy aircraft in ground proximity (110 cases)
- Assessment using rms deviation between measurements and DVM results



# WakeFRA 2004 database

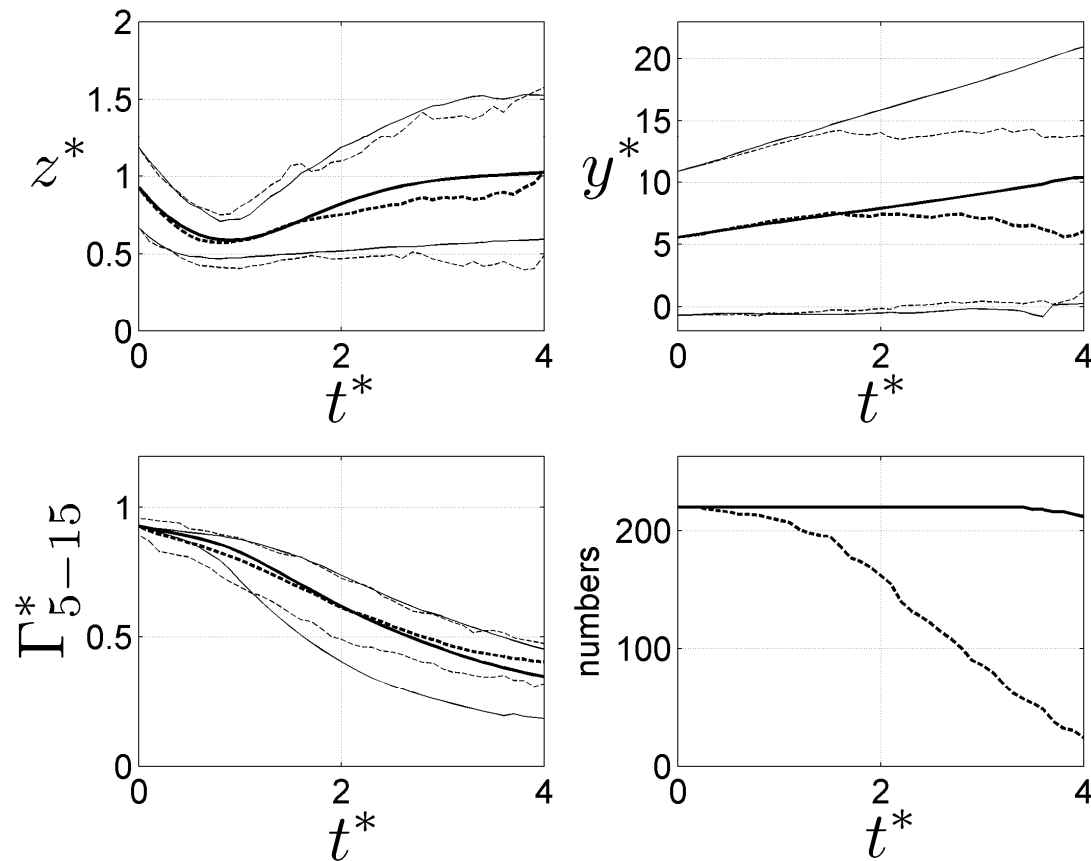
- Results using the FAR-Wake DVM version ("old")



	90 s		110 s	
	old	new	old	new
$\Delta z^*$	0.32	0.26	0.40	0.34
$\Delta y^*$	1.47	1.22	1.72	1.45
$\Delta \Gamma_{5-15}^*$	0.10	0.09	0.09	0.09

# WakeFRA 2004 database

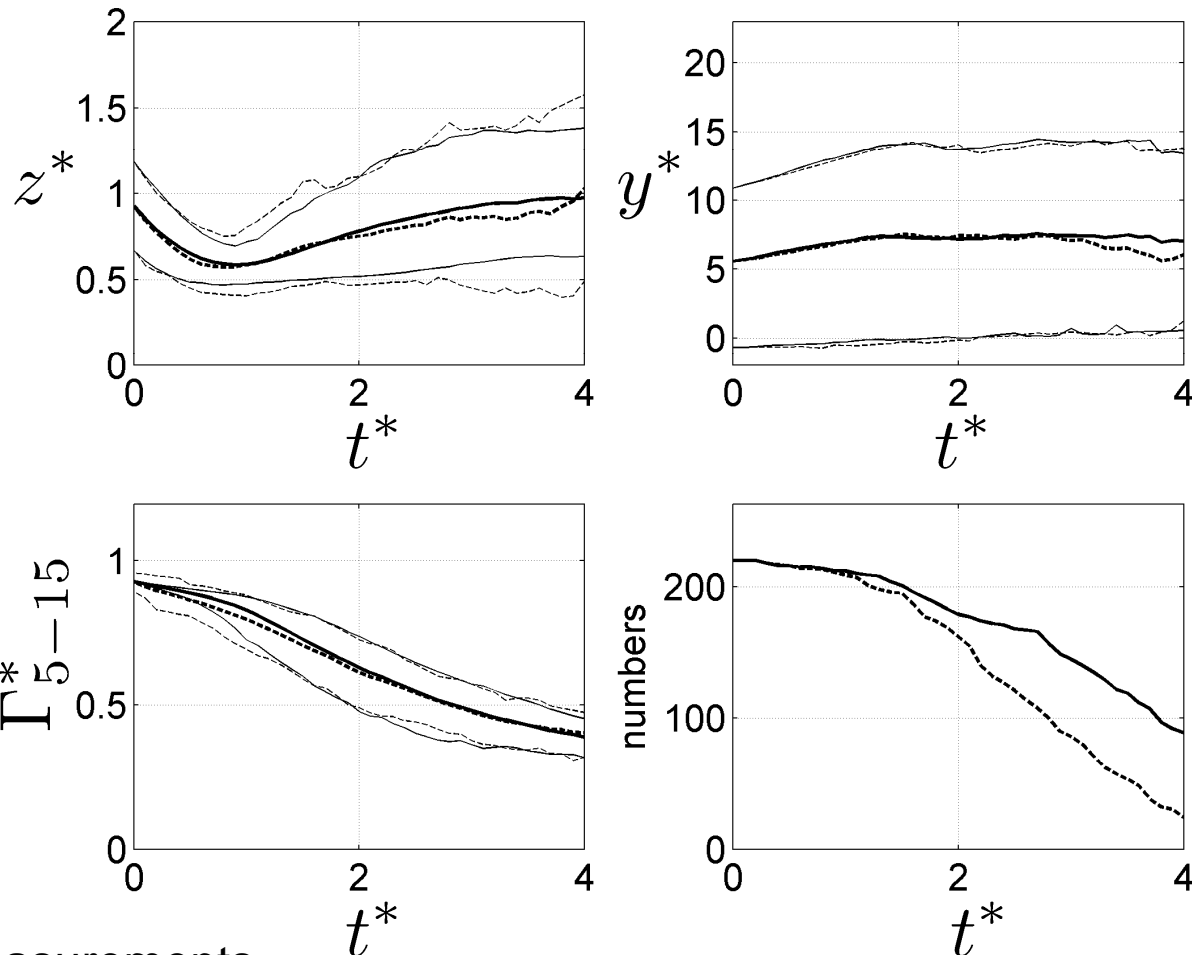
Comparison of the evolution of the mean and 90%-envelope of the whole database



----- Measurements  
—— DVM results

# WakeFRA 2004 database

Use of an equivalent LIDAR-filter on the minimum and maximum measurable lateral position and minimum measurable circulation



----- Measurements  
—— DVM results



# Conclusions on the improved IGE model

- For a case of WV generated IGE, the improved IGE model uses 6 to 9 times less secondary vortex particles.
- It still mimics correctly the flow physics
- Its quality is also improved wrt the previous version.
- It is computationally much more efficient ( $>10$ ).

# Modeling of the effect of stratification combined with weak turbulence

- Stratification affects the transport of the vortices :
  - Rebound of the vortices due to buoyancy
  - Decrease of the vortex lateral separation for moderate  $N^*$
  - Increase of the vortex lateral separation for high  $N^*$
- Stratification also affects the vortex decay
  - Strong interaction with the turbulent baroclinic vorticity
  - For moderate  $N^*$ , increase of the Crow instability growth rate

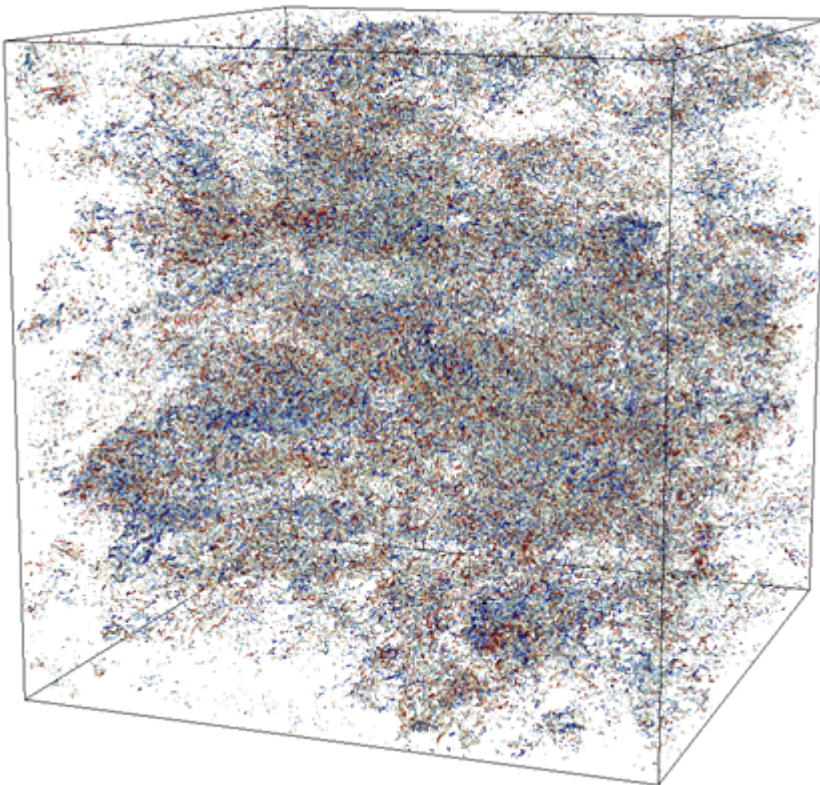
# LES of WV in stratified and weakly turbulent atmosphere

- Pseudo-spectral LES at very high  $Re = \frac{\Gamma_0}{\nu}$  (LES using SGS model only)
- Periodic domain with possible Crow instability development
- Relatively tight vortex core ( $\frac{r_c}{b_0} = 0.05$ ) with fine LES grids ( $\frac{h}{b_0} = \frac{1}{64}$ )
- Five stratification levels are investigated from neutral to very high:  
 $N^* = 0.0, 0.35, 0.75, 1.0, 1.4$
- Two turbulence forcing levels:  $\left(\frac{dE}{dt}\right)_f^* = \frac{dE}{dt} \frac{b_0}{V_0^3} = 2.4 \cdot 10^{-4}$  and  $2.4 \cdot 10^{-5}$

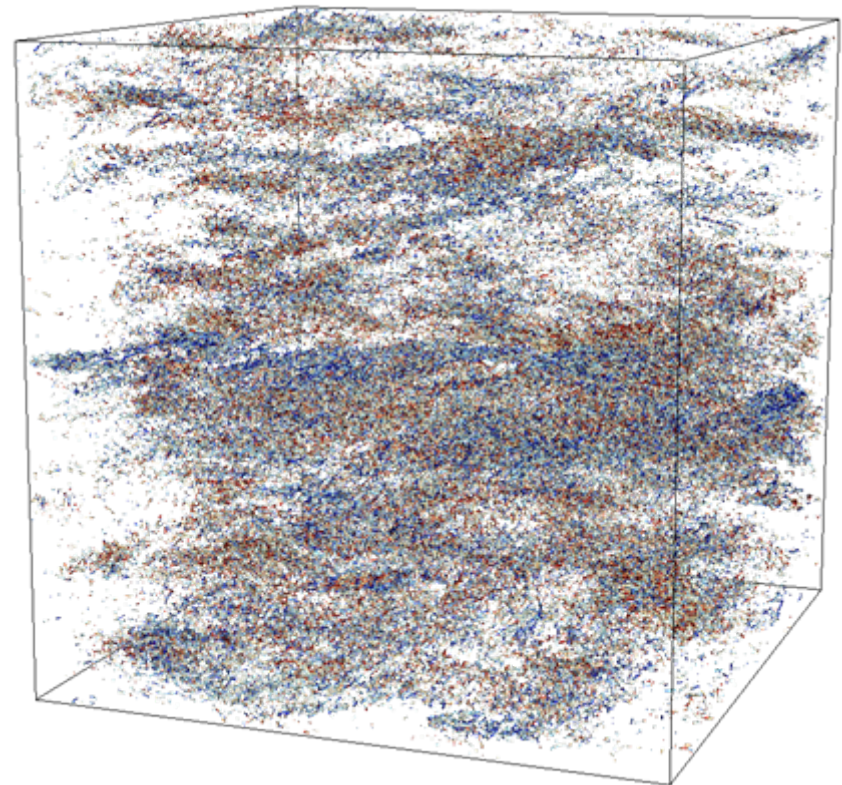
All cases correspond to a low turbulence level attained at equilibrium

# Stratified turbulence

- Pre-simulations of forced stratified turbulence have been performed for the 5 considered  $N^*$ , using LES
- This enables to obtain realistic background stratified and weakly turbulent fields.
- The WV will then evolve in those fields



$N^*=0.35$



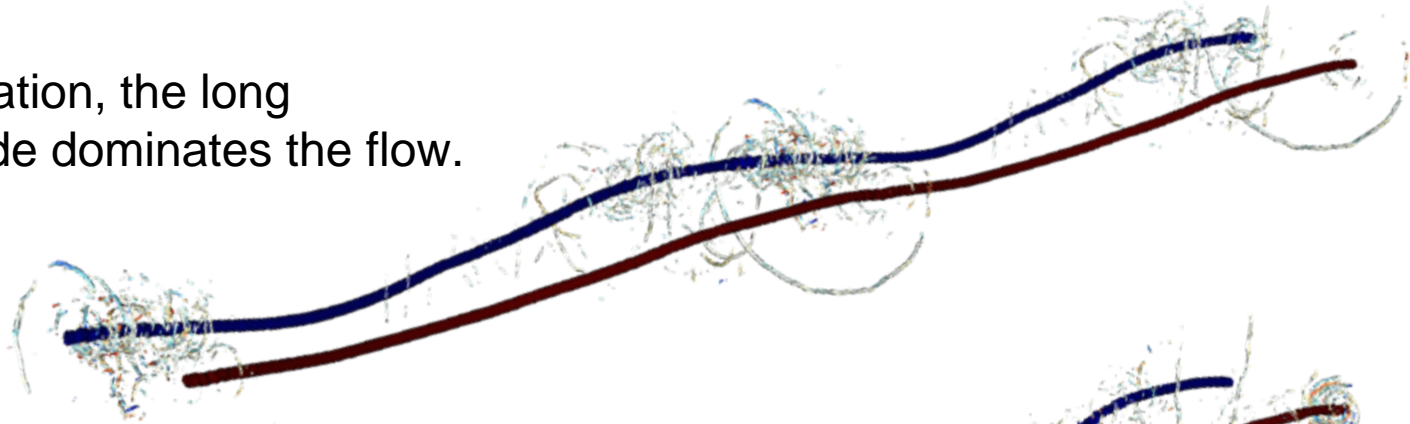
$N^*=1.0$



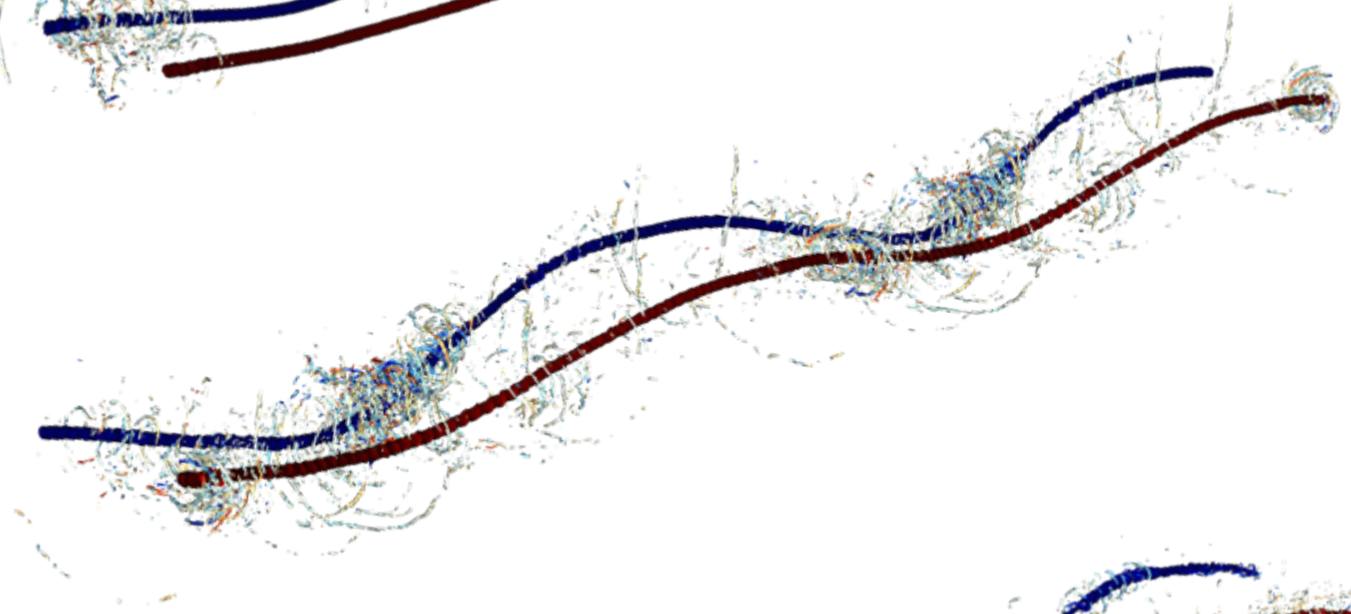
## Case $N^*=0$

Without stratification, the long wavelength mode dominates the flow.

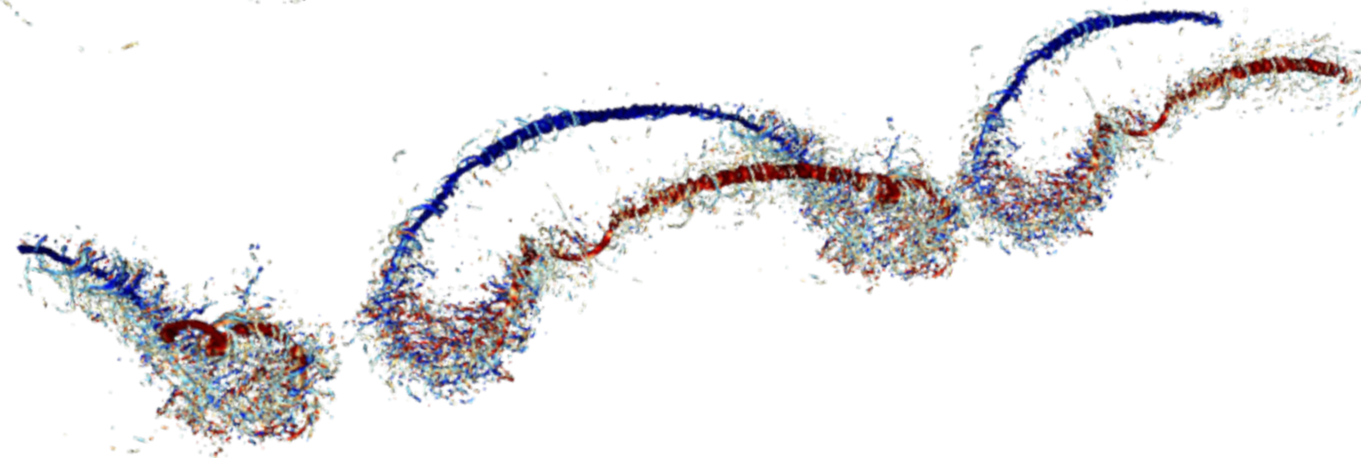
$\tau = 3.0$



$\tau = 4.0$



$\tau = 5.0$

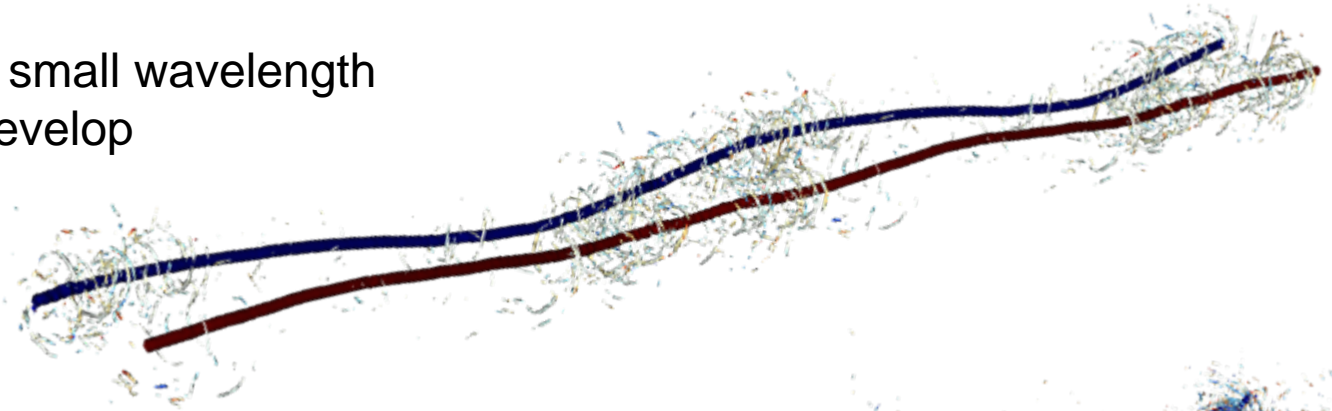


Iso- $\lambda_2$  surfaces colored by the axial vorticity

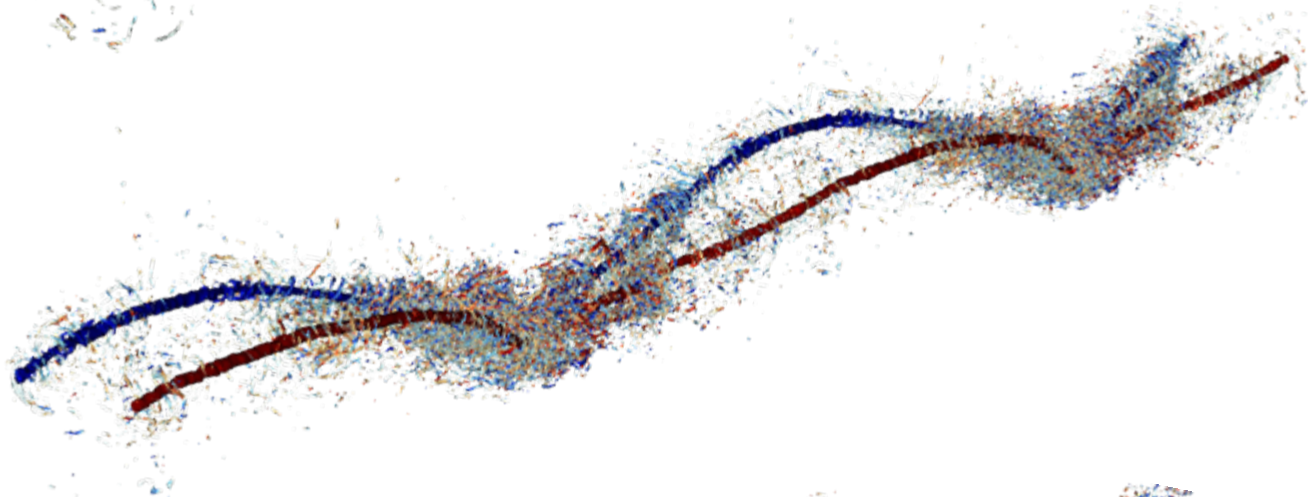
## Case $N^*=0.35$

With stratification, small wavelength instabilities also develop

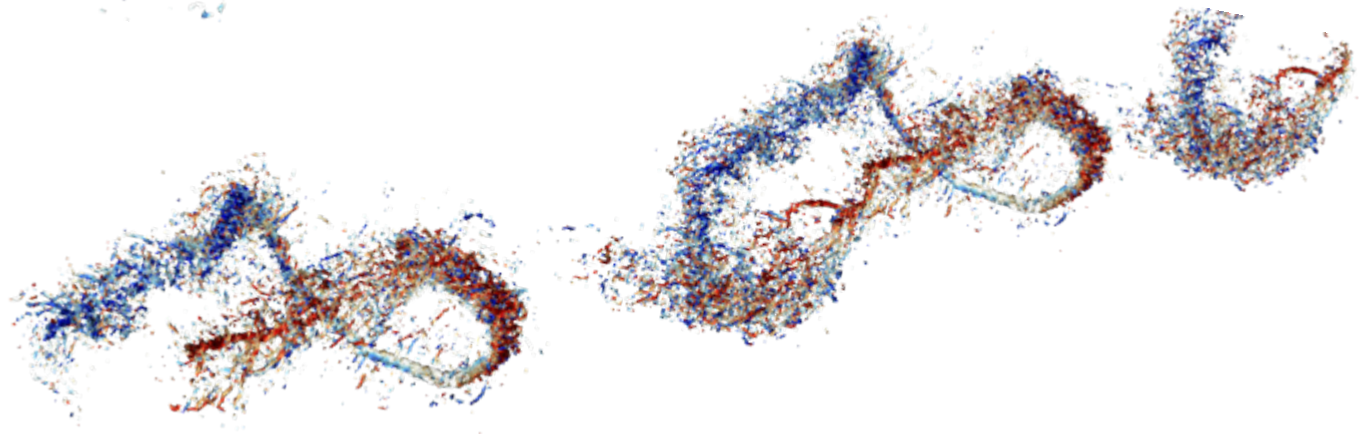
$\tau = 3.0$



$\tau = 4.0$



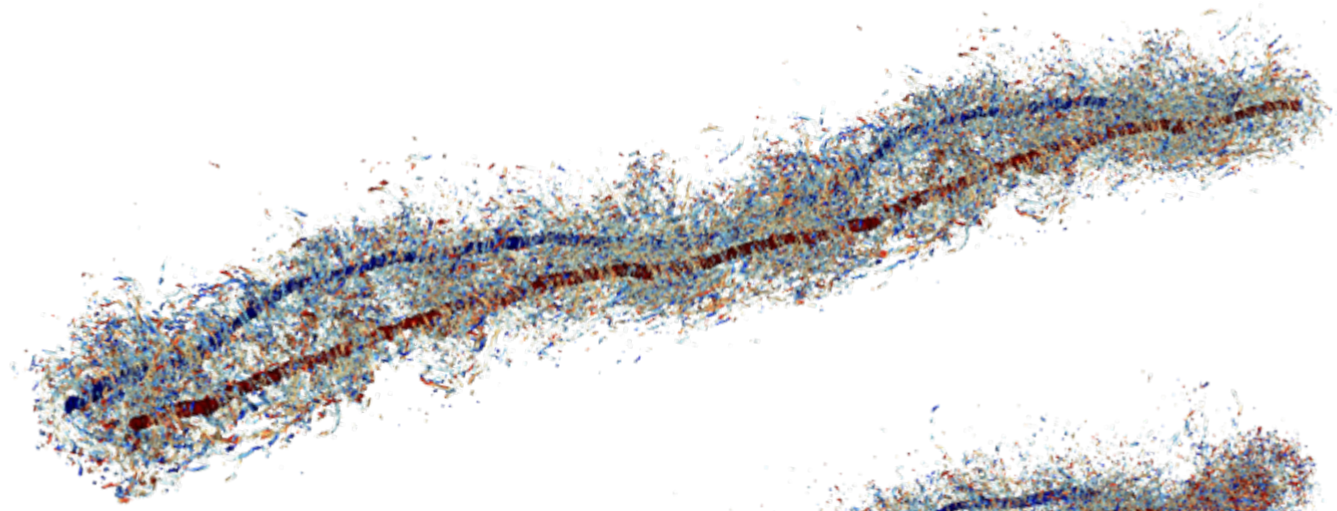
$\tau = 5.0$



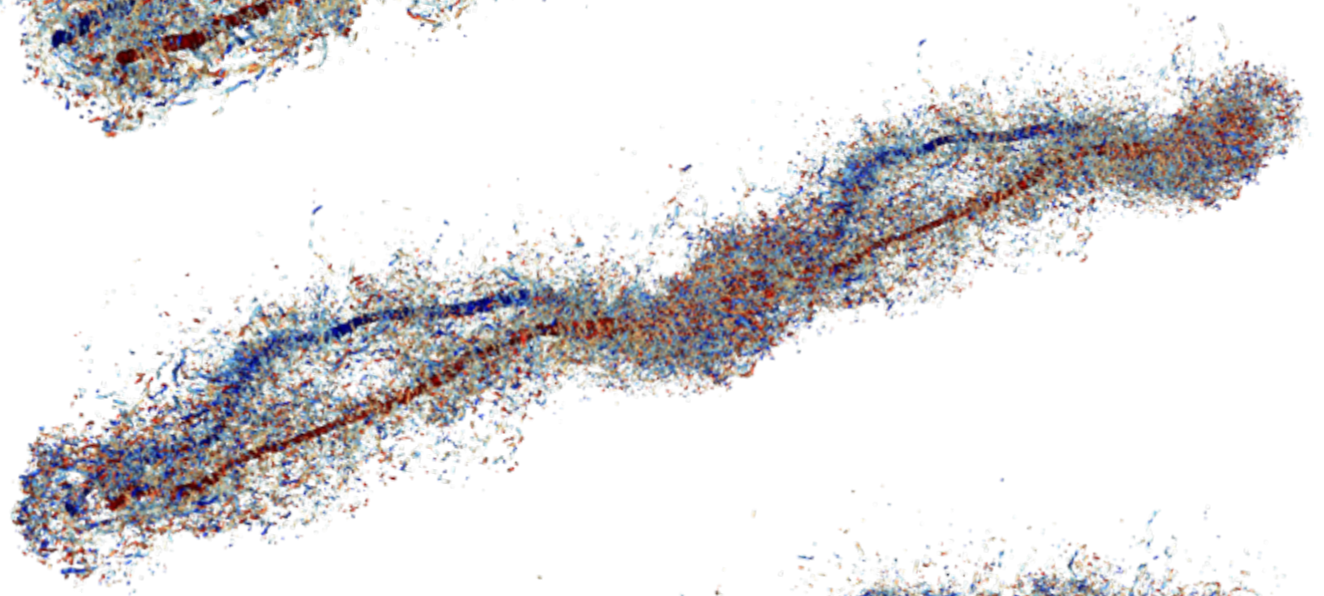
Iso- $\lambda_2$  surfaces colored by the axial vorticity

# Case $N^*=0.75$

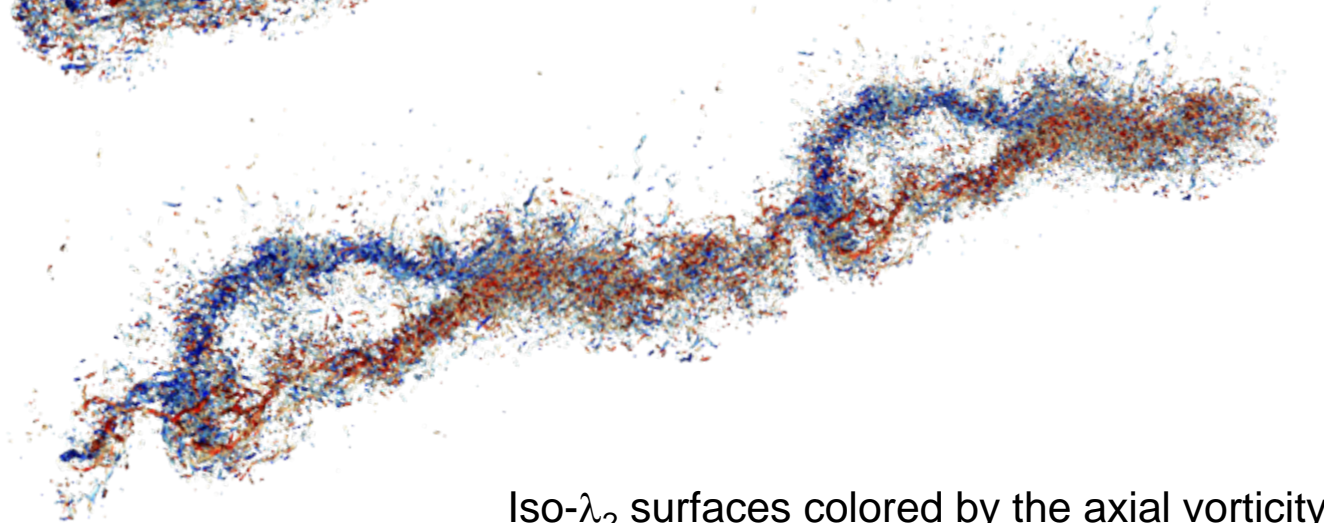
$\tau = 2.5$



$\tau = 3.0$



$\tau = 3.5$



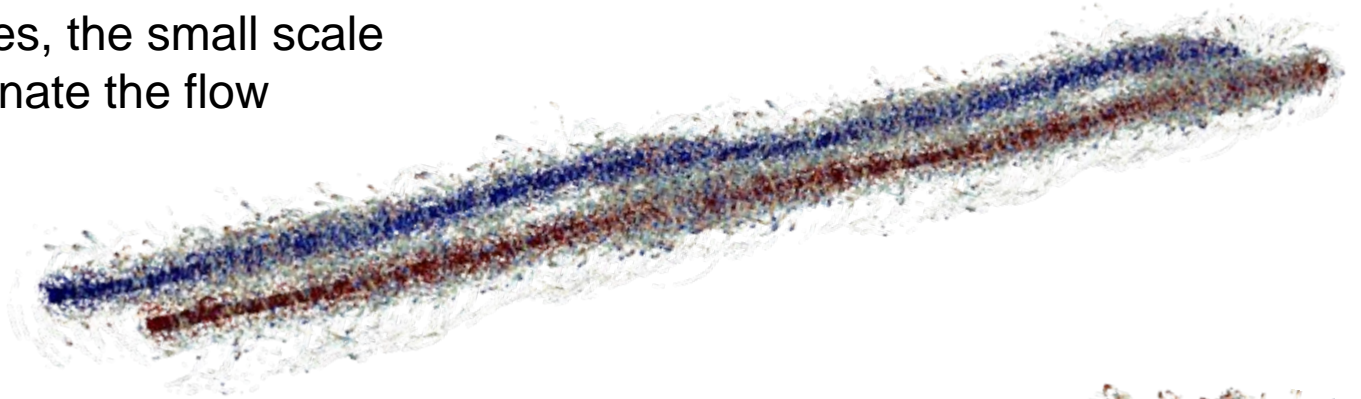
Iso- $\lambda_2$  surfaces colored by the axial vorticity



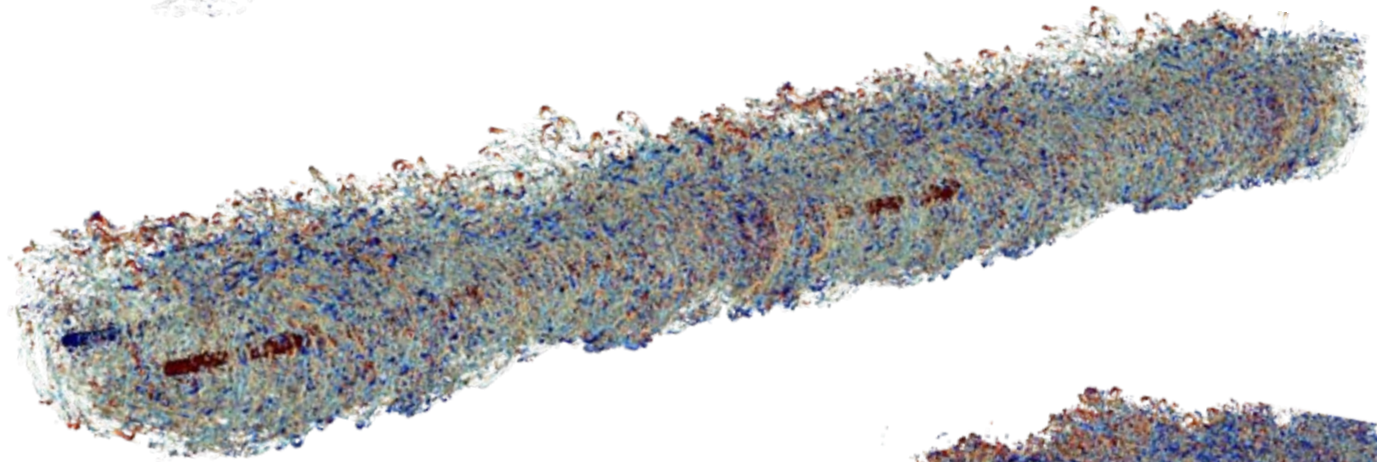
## Case $N^*=1$

For high  $N^*$  values, the small scale instabilities dominate the flow

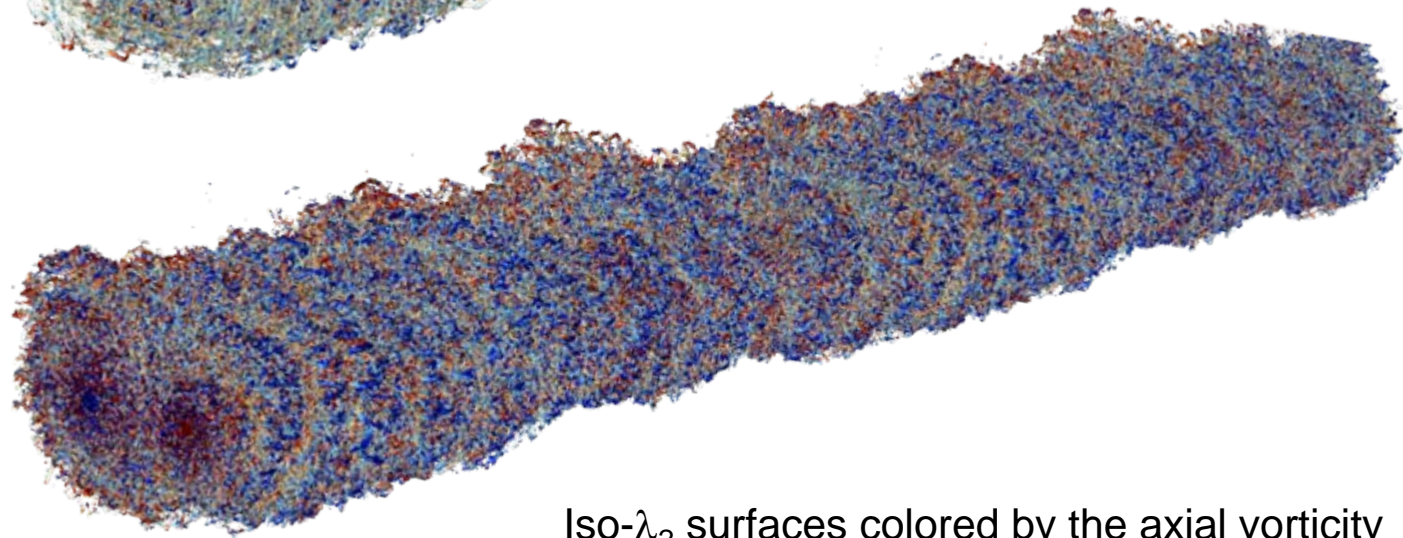
$\tau = 1.0$



$\tau = 1.5$



$\tau = 2.0$



Iso- $\lambda_2$  surfaces colored by the axial vorticity



# Vortex transport modeling

- Vortex altitude evolution influenced by Biot-Savart and stratification :

$$\frac{dz^*}{dt^*} = v_{BS}^* + v_{str}^* \quad \text{with} \quad v_{BS}^*(t^*) = -\Gamma_{tot}^*(t^*)$$

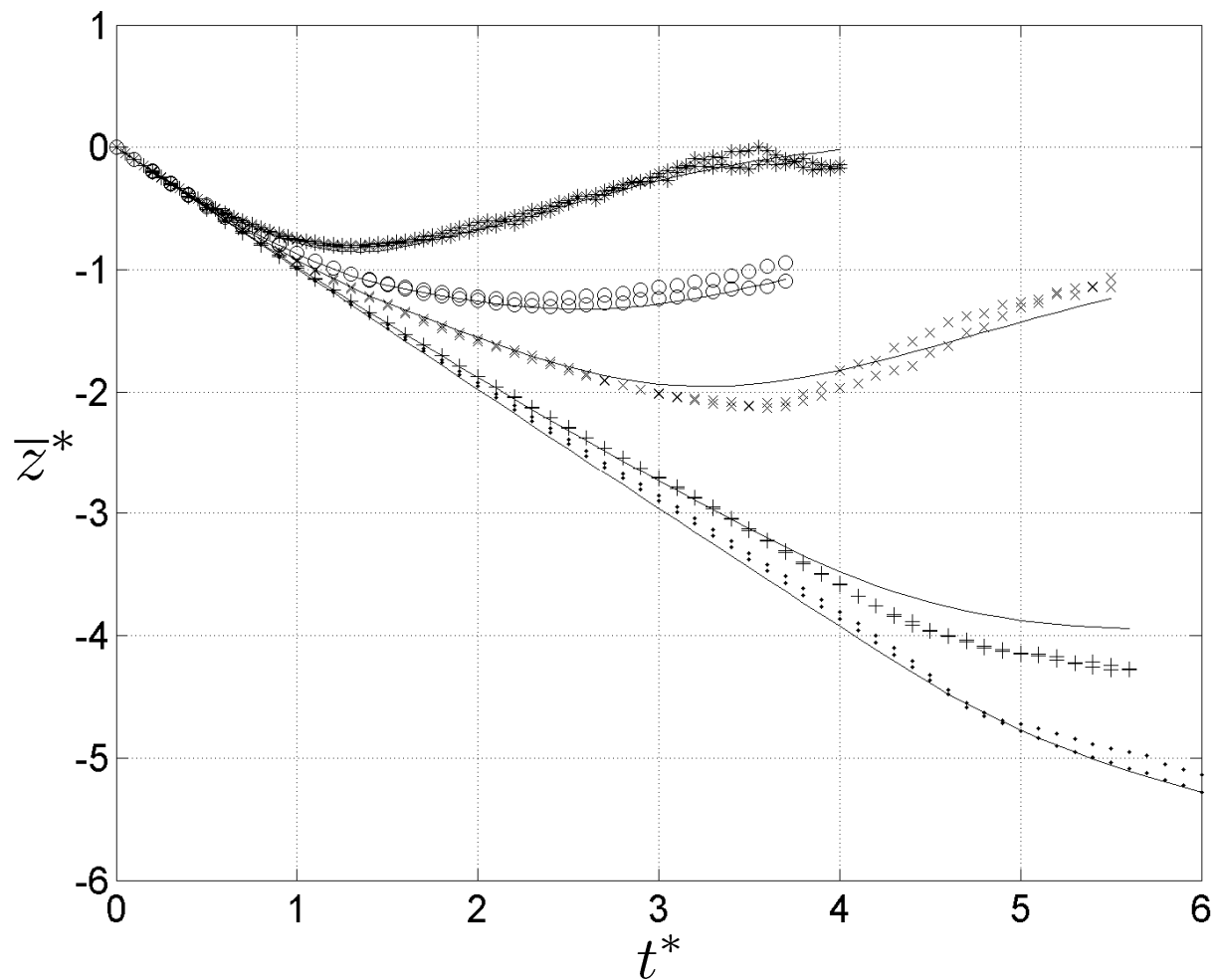
- New model, improved over Saffman's model (1972) :

$$\frac{dv_{str}^*}{dt^*} = - \underbrace{(\alpha_{str} N^*)^2 (z^* - h_0^*)}_{\text{Saffman's model}} - \underbrace{C v_{str}^*}_{\text{Damping term}}$$

- **Saffman's model:** sinusoidal rebound controlled by  $N^*$
- **Damping term :** -modeling the effect of the reorganization of the baroclinic vorticity around the primary vortices, attenuating the rebound effect

$$\left\{ \begin{array}{l} \alpha_{str} = 0.8 \\ C = 1.3 \\ t_{crit}^* = 1.6 \end{array} \right. \quad \text{- Activated for } t^* \geq t_{crit}^*$$

# Comparison with LES data



- $N^* = 0$
- +  $N^* = 0.35$
- ×  $N^* = 0.75$
- $N^* = 1.0$
- \*  $N^* = 1.4$

# Circulation decay modeling

- New time-to-demise model, based on Sarpkaya's model (2000)  $t_{d,S}^*(\epsilon^*)$

$$t_d^*(\epsilon^*, N^*) = \begin{cases} t_{d,S}^* \exp \left( -C_1 \left( N^* t_{d,S}^* \right)^a \right) & \text{if } (N^* t_{d,S}^*) \leq K \\ \frac{C_2}{N^*} & \text{otherwise} \end{cases}$$

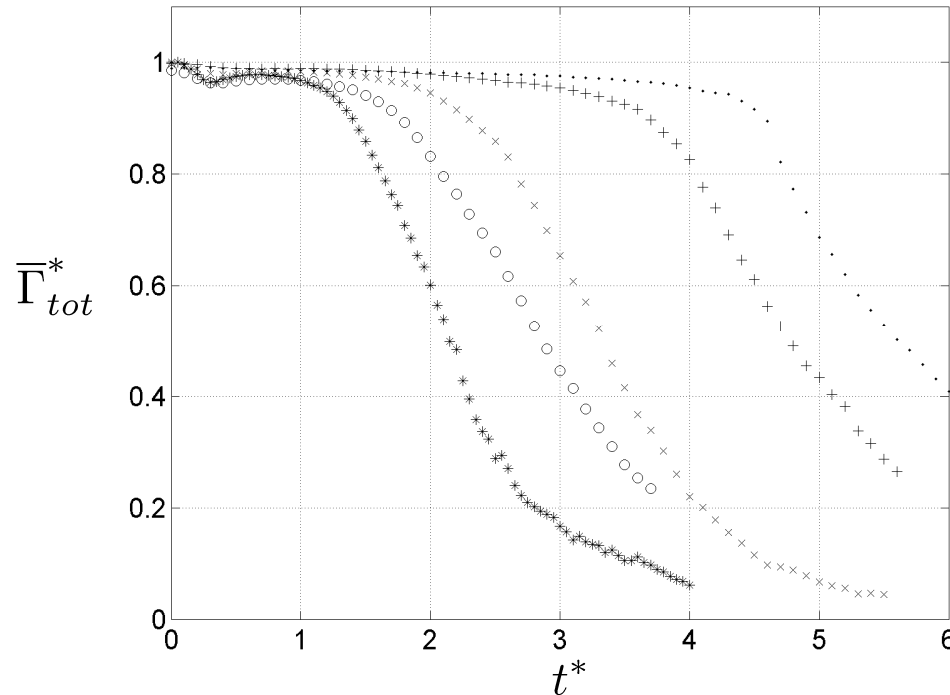
$$\text{with } K = (C_1 a)^{-1/a} \text{ and } C_2 = K \exp(-C_1 K^a)$$

Model established so that:

- $t_d^*$  is a decreasing function of  $\epsilon^*$  and  $N^*$
- for very high  $N^*$ ,  $t_d^*$  is independent on  $\epsilon^*$

$$\begin{cases} a = 1.6 \\ C_1 = 0.076 \end{cases} \Rightarrow \begin{cases} K = 3.73 \\ C_2 = 2.0 \end{cases}$$

# Two-phase decay model

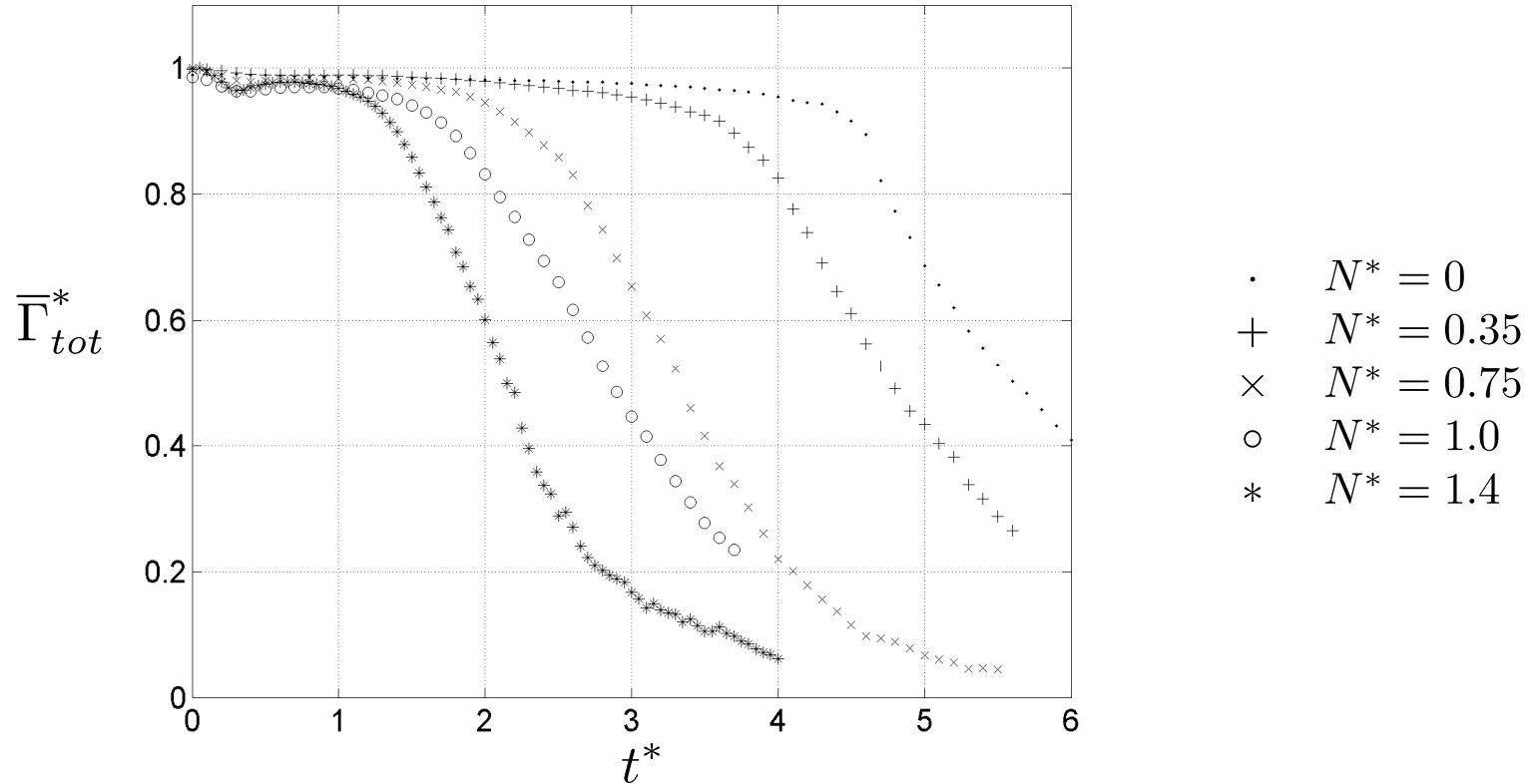


$$\frac{d\Gamma_{tot}^*}{dt^*} = \begin{cases} -\frac{c_1}{t_{d,S}^*} \Gamma_{tot}^* & \text{if } t^* \leq t_d^* \\ -\frac{c_2}{t_{\Gamma}^*} \Gamma_{tot}^* & \text{if } t^* \geq t_d^* \end{cases} \quad \text{with} \quad t_{\Gamma}^* = \begin{cases} t_{d,S}^* & \text{if } N^* = 0, \\ \sqrt{t_{d,S}^* t_N^*} & \text{otherwise.} \end{cases}$$

$$\begin{cases} c_1 & = & 0.035 \\ c_2 & = & 2.75 \\ t_N^* & = & 3.0 \end{cases}$$

# Additional slow decay due to stratification

In the first decay phase, a very slow decay is observed

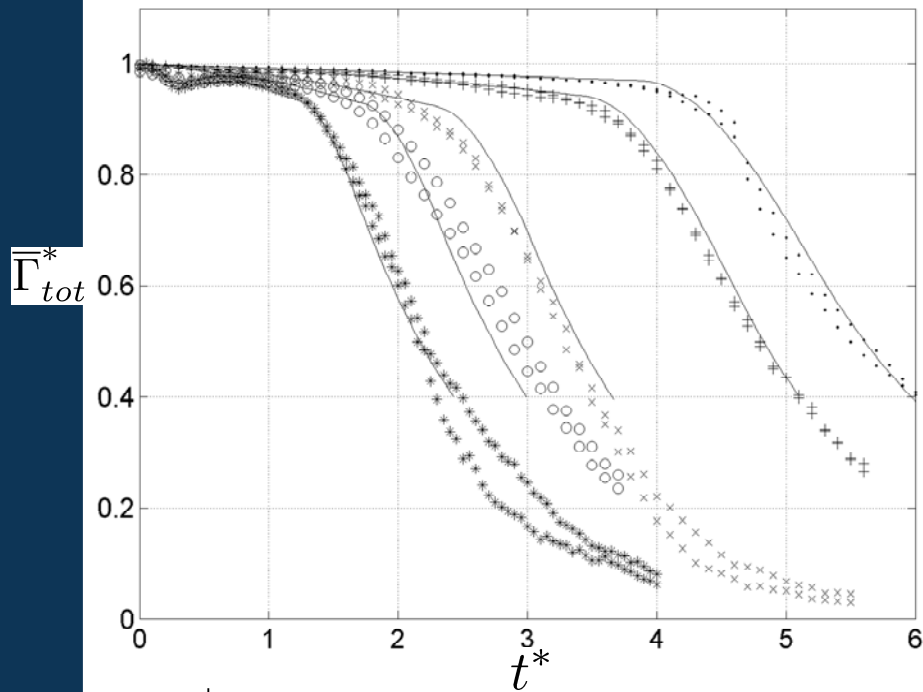


$$\begin{aligned}
 \left. \frac{d\Gamma_{tot}^*}{dt^*} \right|_{str} &= (1 - \beta_{str}) (\alpha_{str} N^*)^2 (z^* - h_0^*) \\
 \frac{dv_{str}^*}{dt^*} &= -\beta_{str} (\alpha_{str} N^*)^2 (z^* - h_0^*) - C v_{str}^*
 \end{aligned}
 \quad \left\{ \begin{array}{l} \beta_{str} = 0.93 \\ \alpha_{str} = 0.83 \end{array} \right.$$

# Comparison with LES data

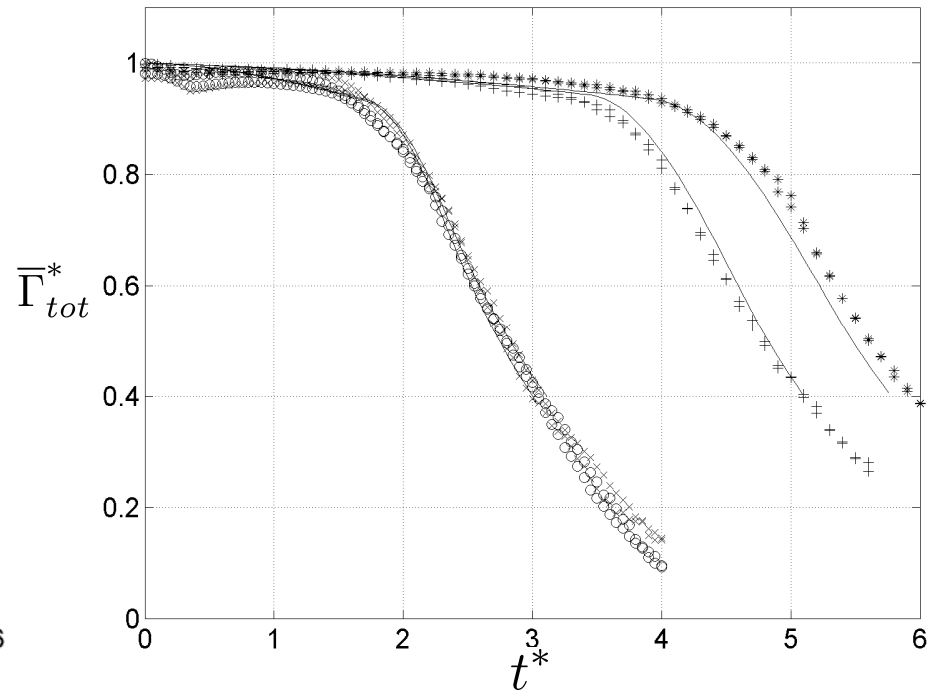
## Longitudinally-averaged circulation evolution

### Stratification effect

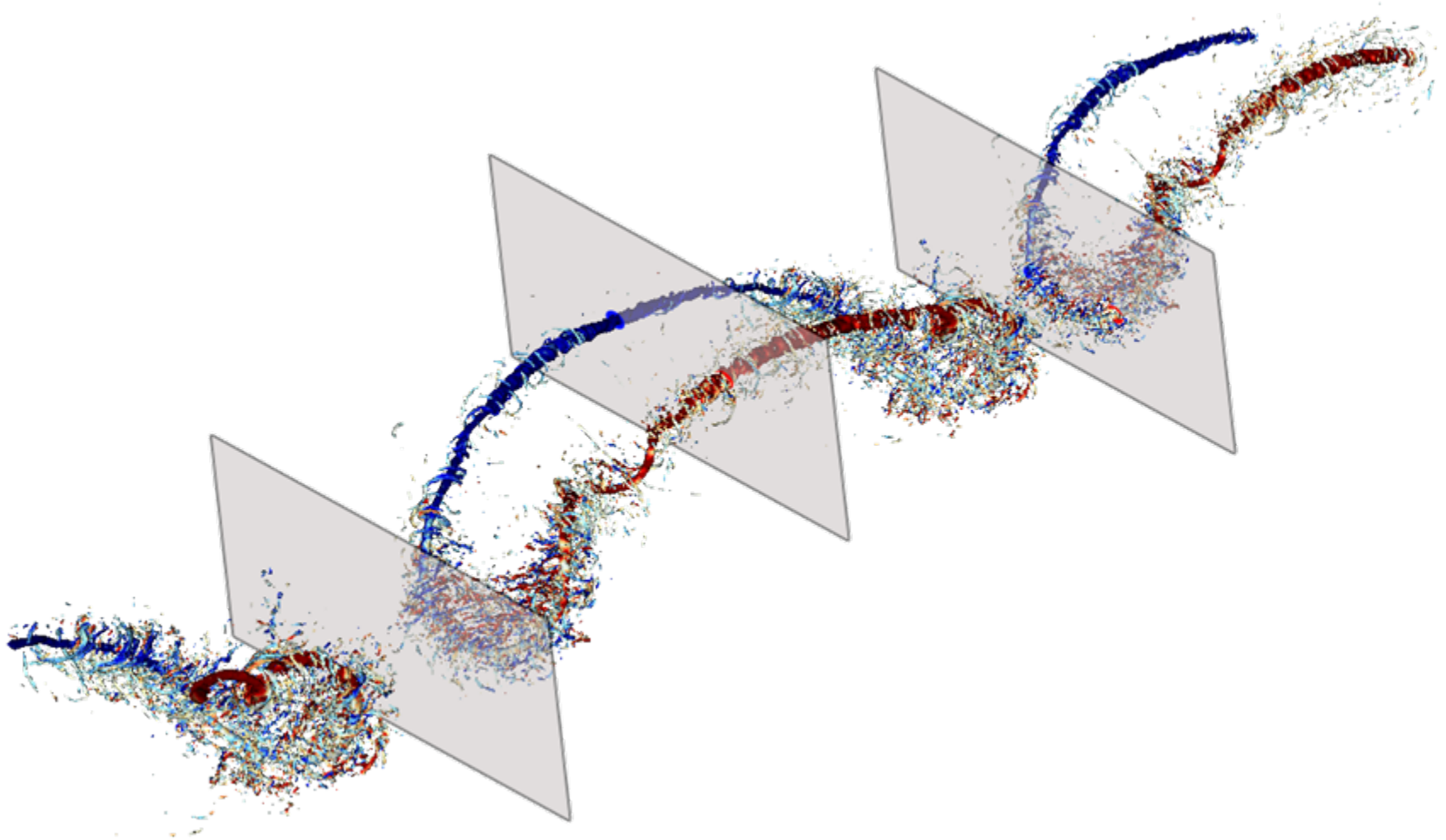


$$\epsilon^* = \frac{\epsilon b_0}{V_0^3} = 2.42 \cdot 10^{-4}$$

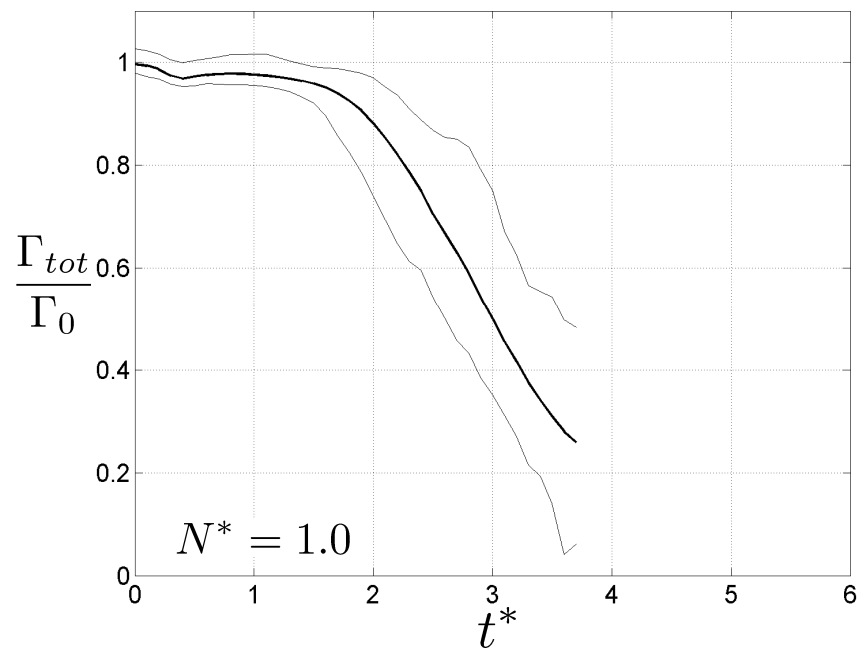
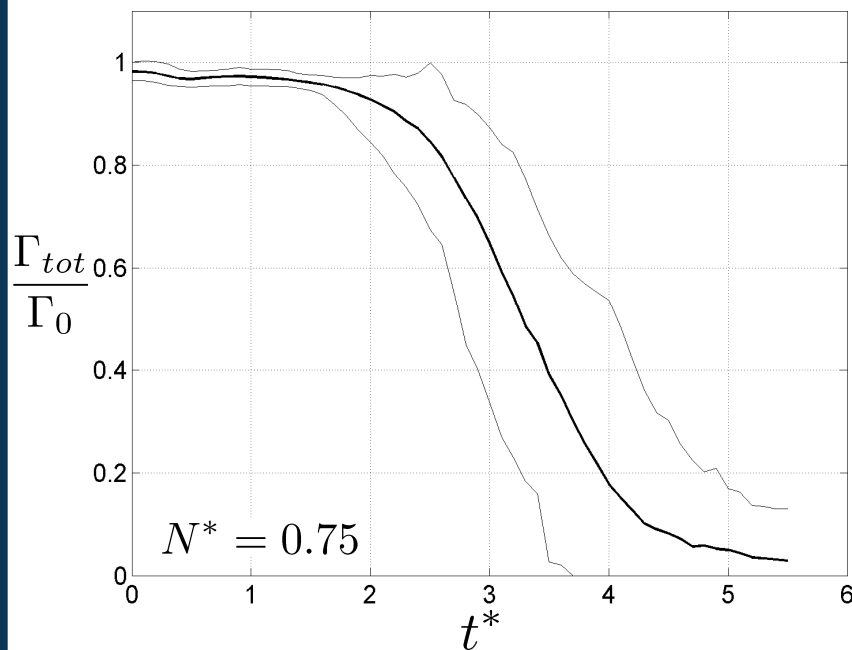
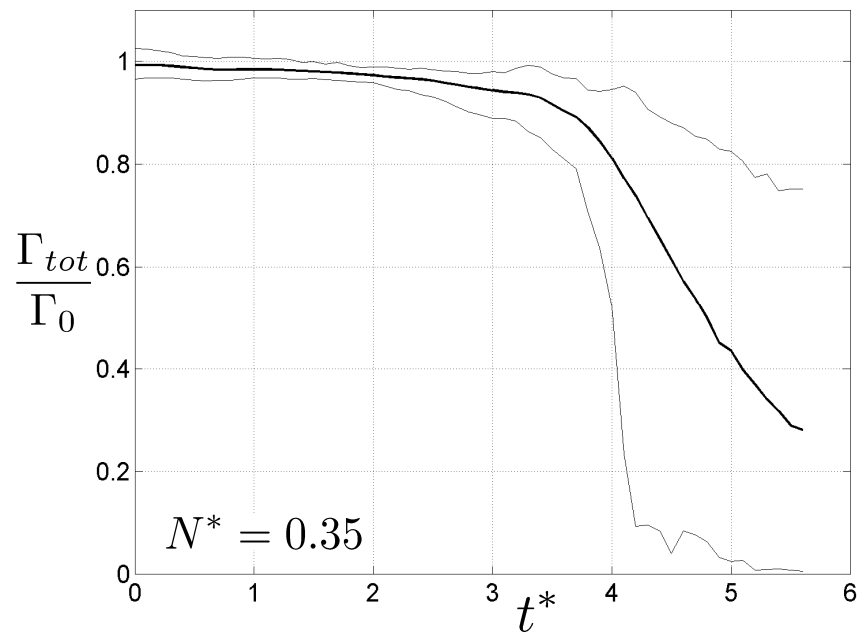
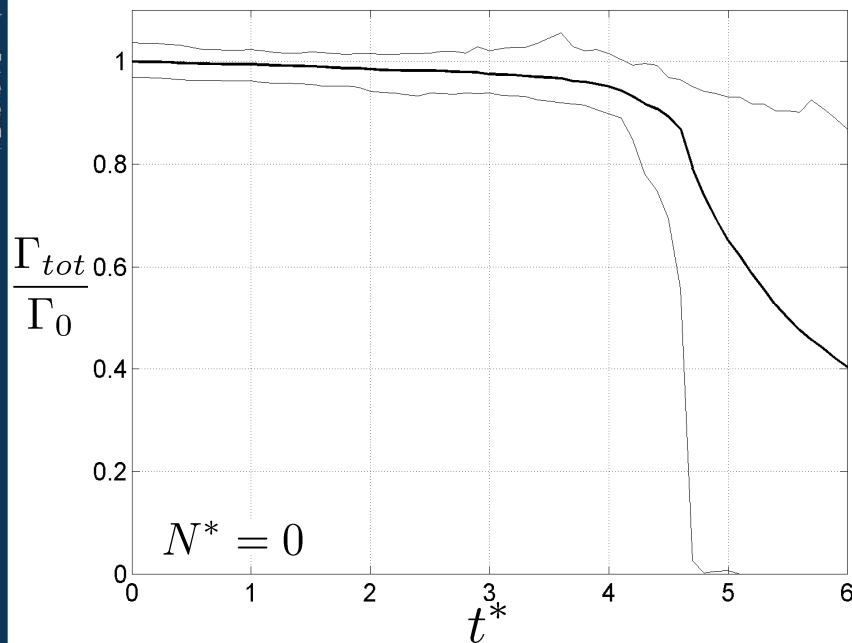
### Stratification and turbulence effects



# Evaluation of the vortex properties in various cross-planes



# Comparison of the mean and 95%-envelope circulation evolution





# Vortex decay

- For low stratification levels:
  - the decay is non uniform along the vortex
  - All planes exhibit a two-phase decay with different fast decay rates
  - Second phase due to Crow linking and stratification
- For high  $N^*$ :
  - The decay is much more uniform
  - Two-phase decay with almost uniform decay rate
- The model coefficients could be adapted, to be used in the PVM, to also predict the envelope behavior
  - => The DVM model coefficient distribution would then be a function of the stratification level

# Conclusions

- An improved IGE model was developed and integrated in the DVM
  - The model uses less particles. Hence, it is more efficient
  - Moreover, the quality of the model was improved
- An improved model of the effects of atmospheric turbulence and stratification was also developed
  - The altitude evolution model correctly reproduces the LES data
  - The new time-to-demise model takes into account the influence of both turbulence and stratification
  - The 2-phase decay model correctly reproduces the LES data
- Those new models are integrated in the DVM/PVM/WAKE4D so that it is more accurate and more efficient
- Moreover, the PVM can now use OPENMP to perform DVM runs in parallel on the different cores of a PC.
  - => gain of another factor of 4 for a PC with 4 cores.